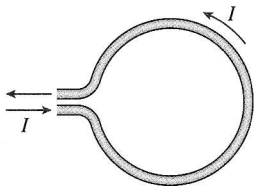


## The magnetic field of a current loop

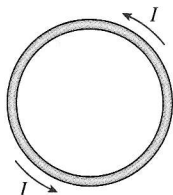
FIGURE 32.16a shows a current loop, a circular loop of wire with radius  $R$  that carries current  $I$ . Find the magnetic field of the current loop at distance  $z$  on the axis of the loop.

FIGURE 32.16 A current loop.

(a) A practical current loop



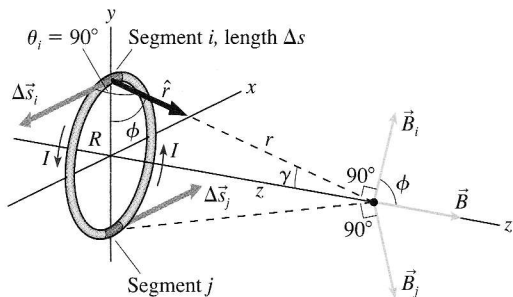
(b) An ideal current loop



**MODEL** Real coils need wires to bring the current in and out, but we'll model the coil as a current moving around the full circle shown in FIGURE 32.16b.

**VISUALIZE** FIGURE 32.17 shows a loop for which we've assumed that the current is circulating ccw. We've chosen a coordinate system in which the loop lies at  $z = 0$  in the  $xy$ -plane. Let segment  $i$  be the segment at the top of the loop. Vector  $\Delta\vec{s}_i$  is parallel to the  $x$ -axis and unit vector  $\hat{r}$  is in the  $yz$ -plane, thus angle  $\theta_i$ , the angle between  $\Delta\vec{s}_i$  and  $\hat{r}$ , is  $90^\circ$ .

FIGURE 32.17 Calculating the magnetic field of a current loop.



The direction of  $\vec{B}_i$ , the magnetic field due to the current in segment  $i$ , is given by the cross product  $\Delta\vec{s}_i \times \hat{r}$ .  $\vec{B}_i$  must be perpendicular to  $\Delta\vec{s}_i$  and perpendicular to  $\hat{r}$ . You should convince yourself that  $\vec{B}_i$  in Figure 32.17 points in the correct direction. Notice that the  $y$ -component of  $\vec{B}_i$  is canceled by the  $y$ -component of magnetic field  $\vec{B}_j$  due to the current segment at the bottom of the loop,  $180^\circ$  away. In fact, every current segment on the loop can be paired with a segment  $180^\circ$  away, on the opposite side of the loop, such that the  $x$ - and  $y$ -components of  $\vec{B}$  cancel and the components of  $\vec{B}$  parallel to the  $z$ -axis add. In other words, the symmetry of the loop requires the on-axis magnetic field to point along the  $z$ -axis. Knowing that we need to sum only the  $z$ -components will simplify our calculation.

**SOLVE** We can use the Biot-Savart law to find the  $z$ -component  $(B_i)_z = B_i \cos \phi$  of the magnetic field of segment  $i$ . The cross product  $\Delta\vec{s}_i \times \hat{r}$  has magnitude  $(\Delta s)(1) \sin 90^\circ = \Delta s$ , thus

$$(B_i)_z = \frac{\mu_0 I \Delta s}{4\pi r^2} \cos \phi = \frac{\mu_0 I \cos \phi}{4\pi(z^2 + R^2)} \Delta s$$

where we used  $r = (z^2 + R^2)^{1/2}$ . You can see, because  $\phi + \gamma = 90^\circ$ , that angle  $\phi$  is also the angle between  $\hat{r}$  and the radius of the loop. Hence  $\cos \phi = R/r$ , and

$$(B_i)_z = \frac{\mu_0 IR}{4\pi(z^2 + R^2)^{3/2}} \Delta s$$

The final step is to sum the magnetic fields due to all the segments:

$$B_{\text{loop}} = \sum_i (B_i)_z = \frac{\mu_0 IR}{4\pi(z^2 + R^2)^{3/2}} \sum_i \Delta s$$

In this case, unlike the straight wire, none of the terms multiplying  $\Delta s$  depends on the position of segment  $i$ , so all these terms can be factored out of the summation. We're left with a summation that adds up the lengths of all the small segments. But this is just the total length of the wire, which is the circumference  $2\pi R$ . Thus the on-axis magnetic field of a current loop is

$$B_{\text{loop}} = \frac{\mu_0 IR}{4\pi(z^2 + R^2)^{3/2}} 2\pi R = \frac{\mu_0}{2} \frac{IR^2}{(z^2 + R^2)^{3/2}}$$