

*How far away from their black holes  
are quasar outflows located?*

Dunn et al. 2010 (ApJ 709, 611)

Arav et al. 2008 (ApJ 681, 954)

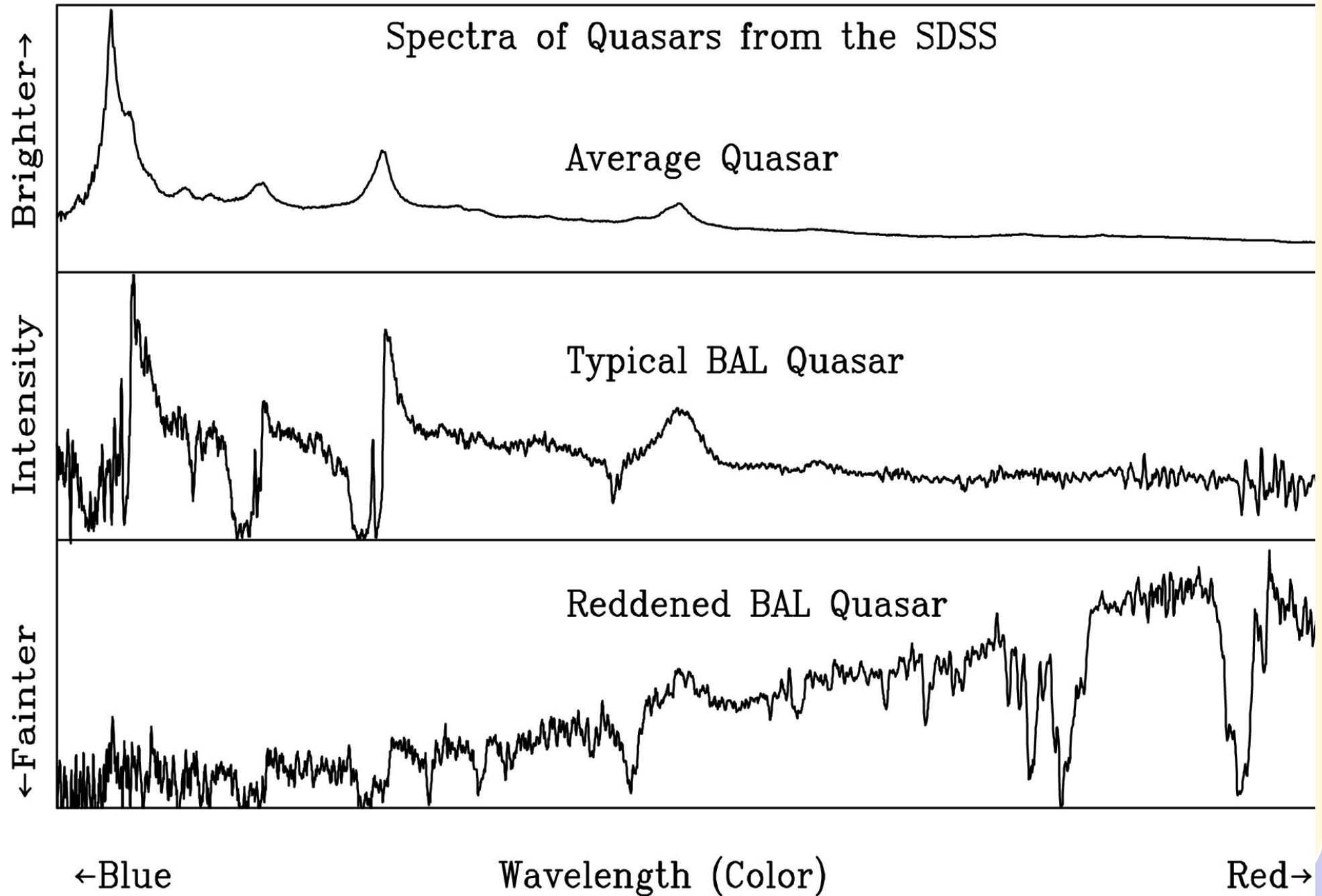
[Korista+ 2008, Moe+ 2009, Bautista+ 2010]

[Chajet et al. 2011]

# Feedback from Quasar Outflows

- Some 20% of quasars show broad absorption line (BAL) troughs;  $f_{BAL} \simeq 0.2$
- Outflow velocities from 0 to 60,000 km/s
- Velocity widths  $> 1000$  km/s (mini-BALs) or  $> 2000$  km/s
- Most common BAL trough is C IV 1548,1550 doublet ( $\Delta v = 500$  km/s); always blended, almost always saturated
- Less common are Mg II 2798,2803 (770 km/s) and Fe II (many!)
- How important are BAL outflows as feedback mechanisms?  
Want to know mass-loss rate, kinetic luminosity, momentum flux.

SDSS targets included normal quasars, 'typical' BAL quasars...



# Determining the Mass-Loss Rate

- Assume outflow has mass  $m$  in thin shell at radius  $R$  which covers fractional solid angle  $\Omega$  as seen from the quasar. Then  $m = 4\pi R^2 \Omega \mu m_p N_H$  where  $\mu m_p$  is the mean mass per particle and  $N_H = \int n_H dR$  is the total hydrogen column density along our sightline.
- Spectra give us the outflow's velocity  $v$  along our sightline.
- Minimum avg. mass loss rate: assume mass  $m$  ejected time  $t = R/v$  ago into fixed  $\Omega$ . Then (assuming one trough only)  $\dot{M}_{min} = m/t = 4\pi \mu m_p R \Omega N_H v$
- Measure  $N_{ion}$ ; need  $N_H$  and  $R$  to get  $\dot{M}_{min}$ .
- Can constrain  $\Omega \leq f_{BAL}$  (due to obscuration).

# Determining the Kinetic Luminosity

- To find  $N_H$  and  $R$ , first find  $n_e \simeq n_H$  using collisionally excited to ground state population ratios of C II, Si II, Fe II, Ni II...
- Next, model the ionization structure of a constant-density slab with ionization parameter  $U_H$  at its face:

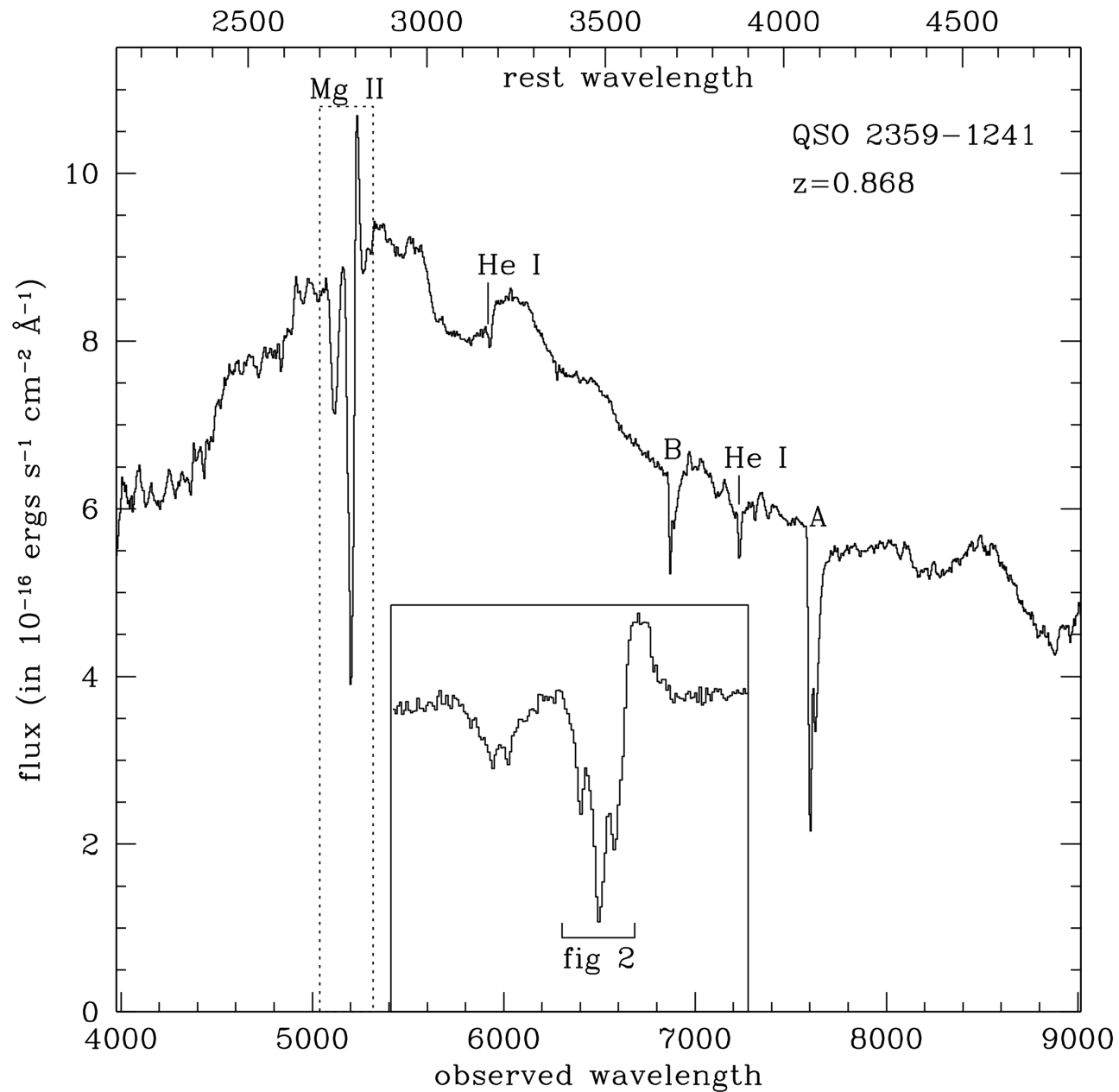
$$U_H = Q_H / 4\pi R^2 c n_H$$

where  $Q_H$  is the # of  $H$ -ionizing photons  $\text{s}^{-1}$  from the quasar:

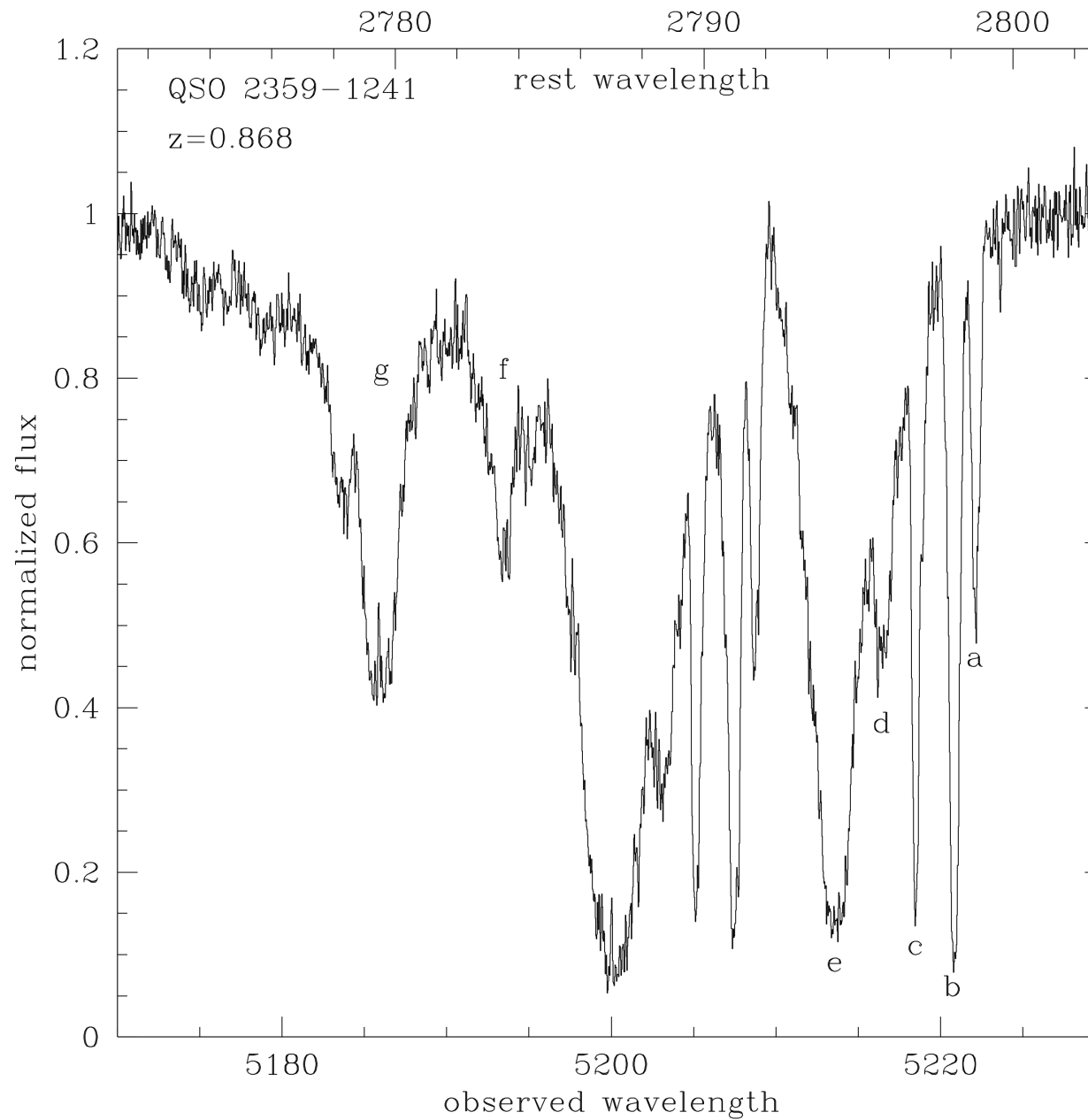
$$Q_H = \int_{1 Ry}^{\infty} \frac{L_\nu}{h\nu} d\nu$$

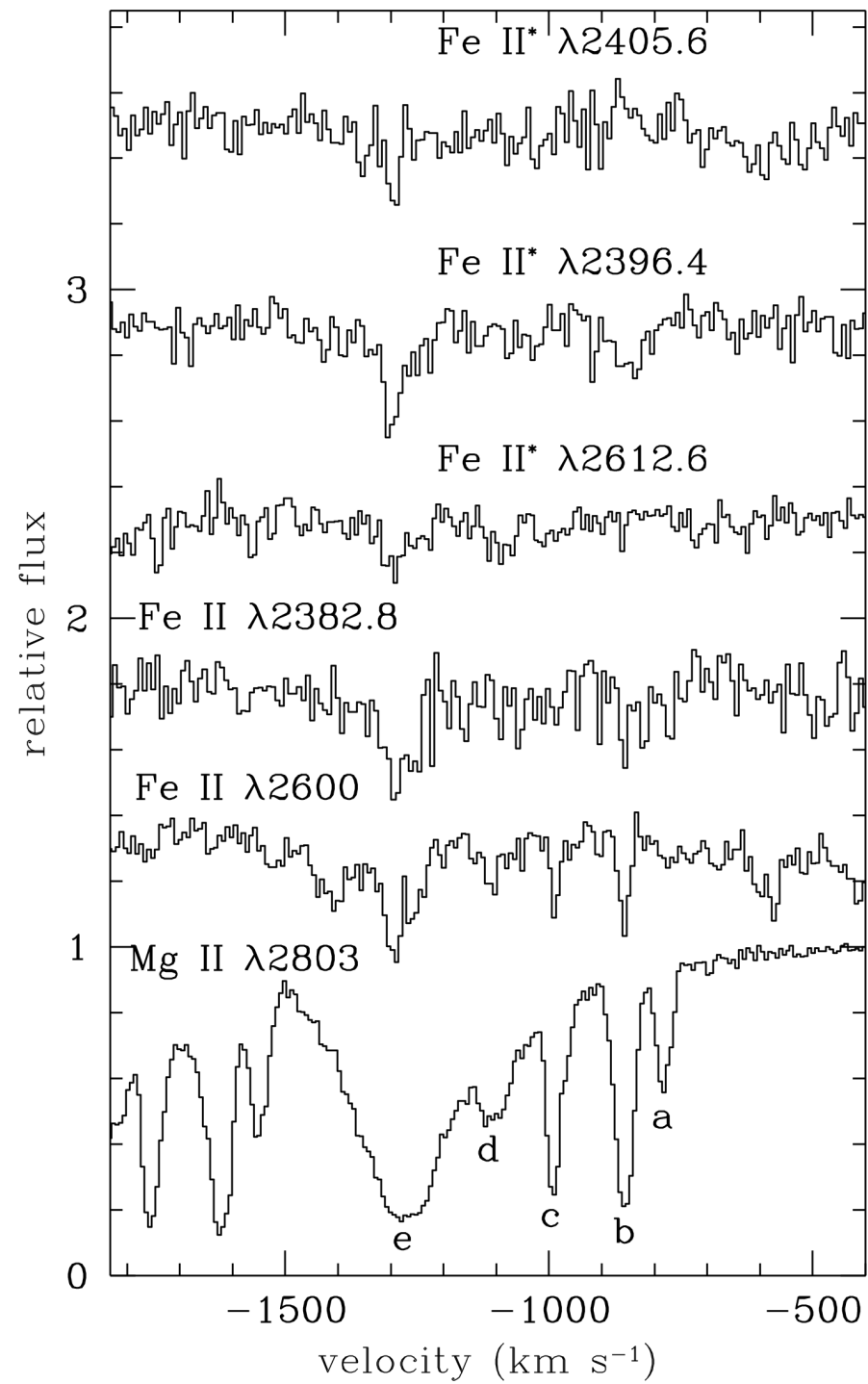
- Find value of  $U_H$  and column density  $N_H$  at which predicted column densities of observed ions best match observations. Need to explore ranges of plausible  $L_\nu$  and metallicity to find best fit and uncertainties for  $N_H$  and  $R$ .
- Kinetic Luminosity is  $\dot{E}_k = \frac{1}{2}\dot{m}v^2 = 2\pi\mu m_p \Omega R N_H v^3$

# Low-resolution spectrum of Arav et al. 2008 quasar



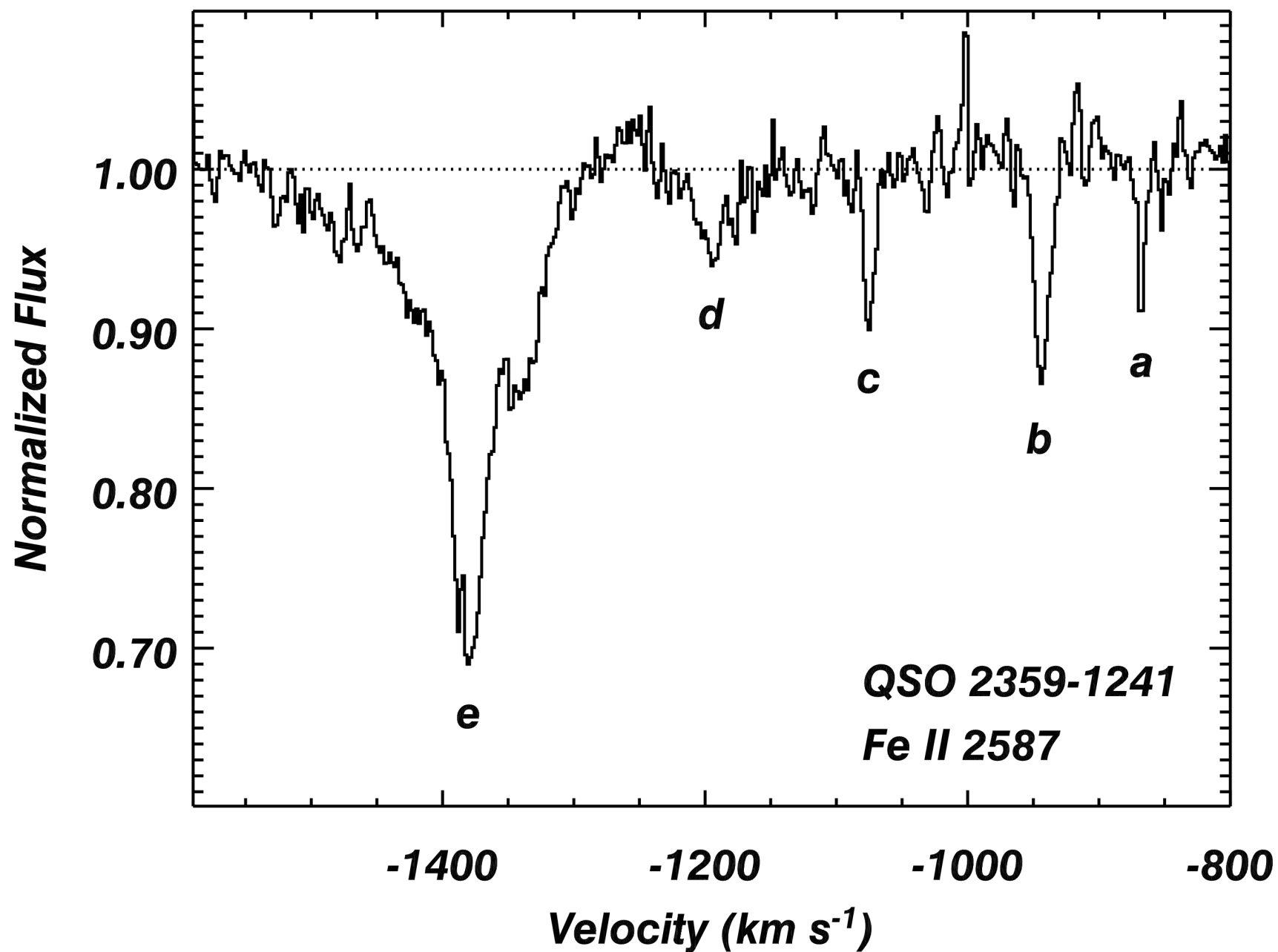
# High-resolution spectrum of Mg II region







# One Fe II transition (2010 data)



## From observed residual intensities to physical parameters:

- Imagine an absorber of optical depth  $\tau$  in some transition in front of a background source with intensity  $I_{\lambda}^{src}$ .
- Complete covering:  $I_{\lambda}^{out} = I_{\lambda}^{src} e^{-\tau}$  or  $I_{\lambda} \equiv I_{\lambda}^{out} / I_{\lambda}^{src} = e^{-\tau}$ .

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**Define the covering factor  $C \leq 1$  so that  $I_{\lambda} = 1 - C + C e^{-\tau}$ .**

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- For doublets, 2 equations & 2 unknowns  $[C(v), \tau(v)]$  at each  $v$ : can always get a solution.
- **When  $\geq 3$  transitions from same ion are available, can check how good an approximation partial covering is.**
- **In many cases, a better approximation is needed.**

How best to approximate complex absorbing structures?





## From observed residual intensities to physical parameters:

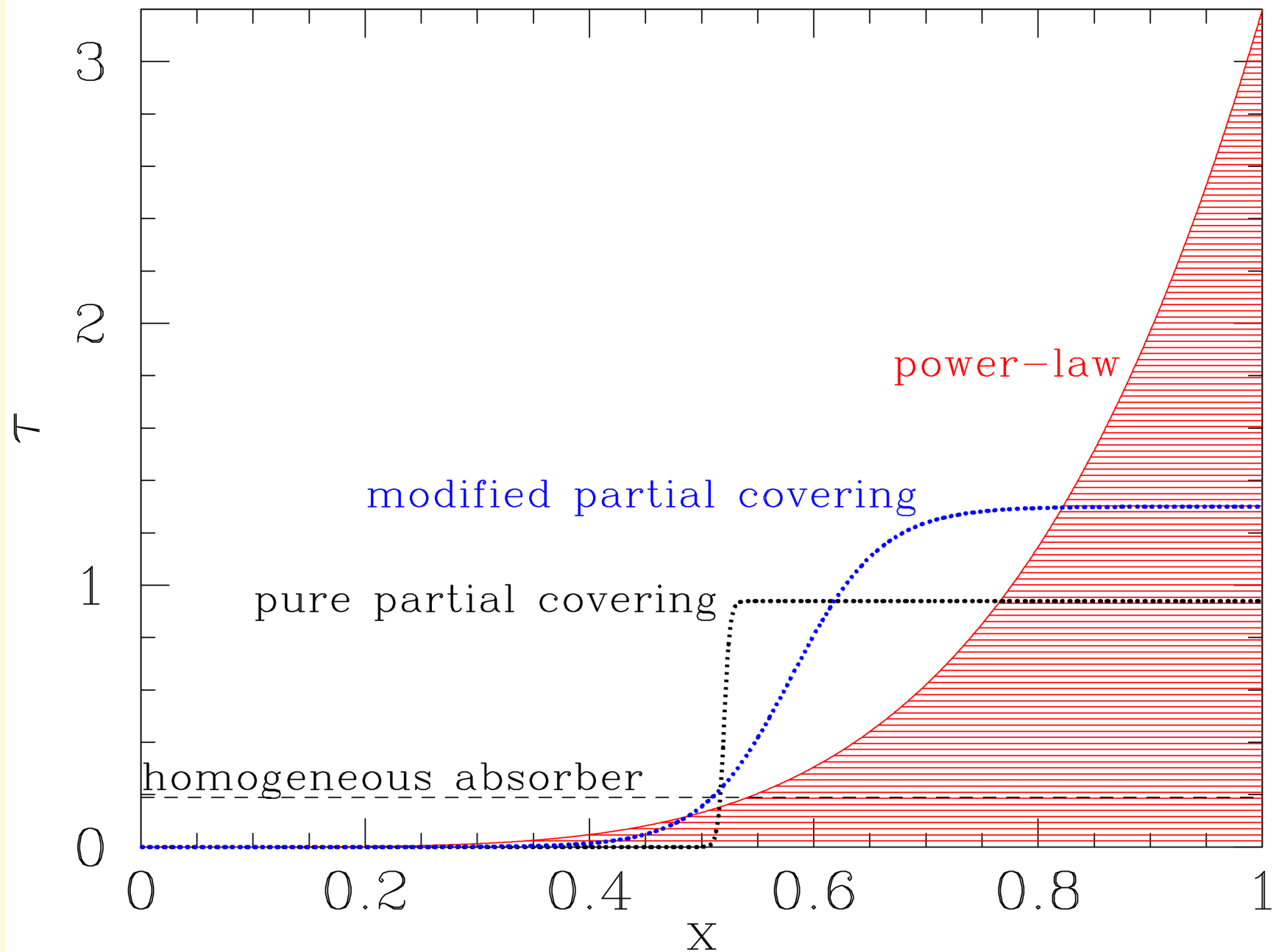
- **Complete covering:**  $I_{\lambda}^{out} = I_{\lambda}^{src} e^{-\tau}$  **or**  $I_{\lambda} \equiv I_{\lambda}^{out} / I_{\lambda}^{src} = e^{-\tau}$
- **Partial covering:**  $I_{\lambda}(v) = 1 - C(v)[1 - e^{-\tau(v)}]$
- **General inhomogeneous absorber:**  
$$e^{-\tau(v)} = \int_x \int_y e^{-\tau(x,y)} dx dy$$



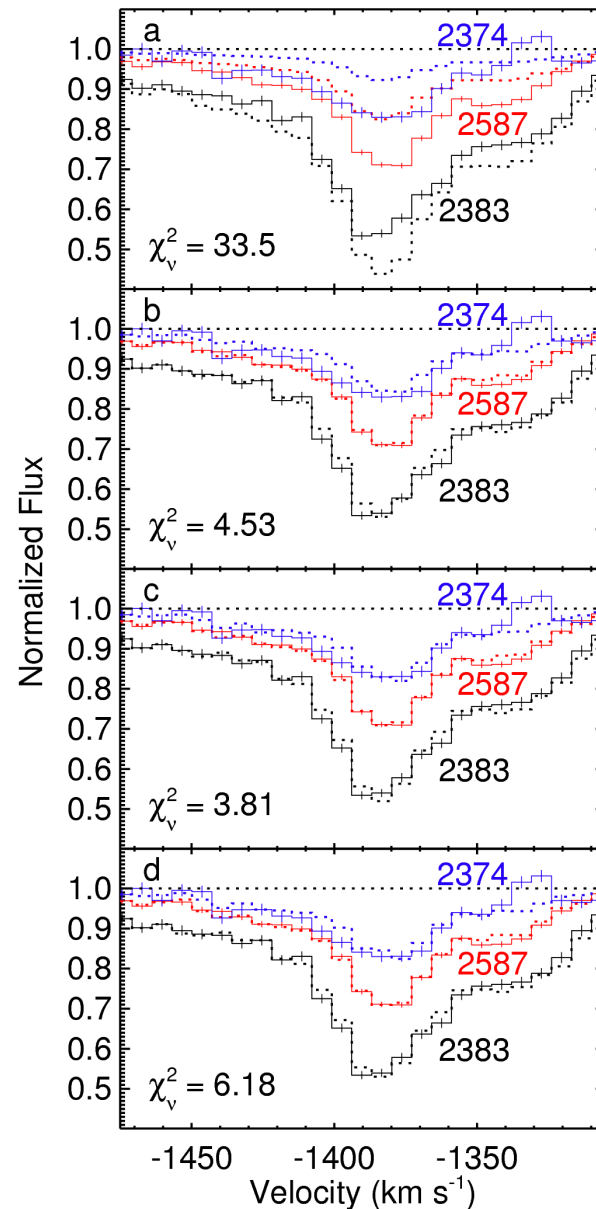
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- **General inhomogeneous absorber: collapse to one dimension, and adopt power-law distribution of optical depths**  
 $\tau(v) = \int_x \tau_{max} x^a dx :$   
 $I_{\lambda}(v) = \int_x \exp(-\tau_{max}(v) x^a) dx$
- **Alternatively, modify partial covering by adding 3rd parameter (width of transition from  $\tau = 0$  to  $\tau = \tau_{max}$ )**

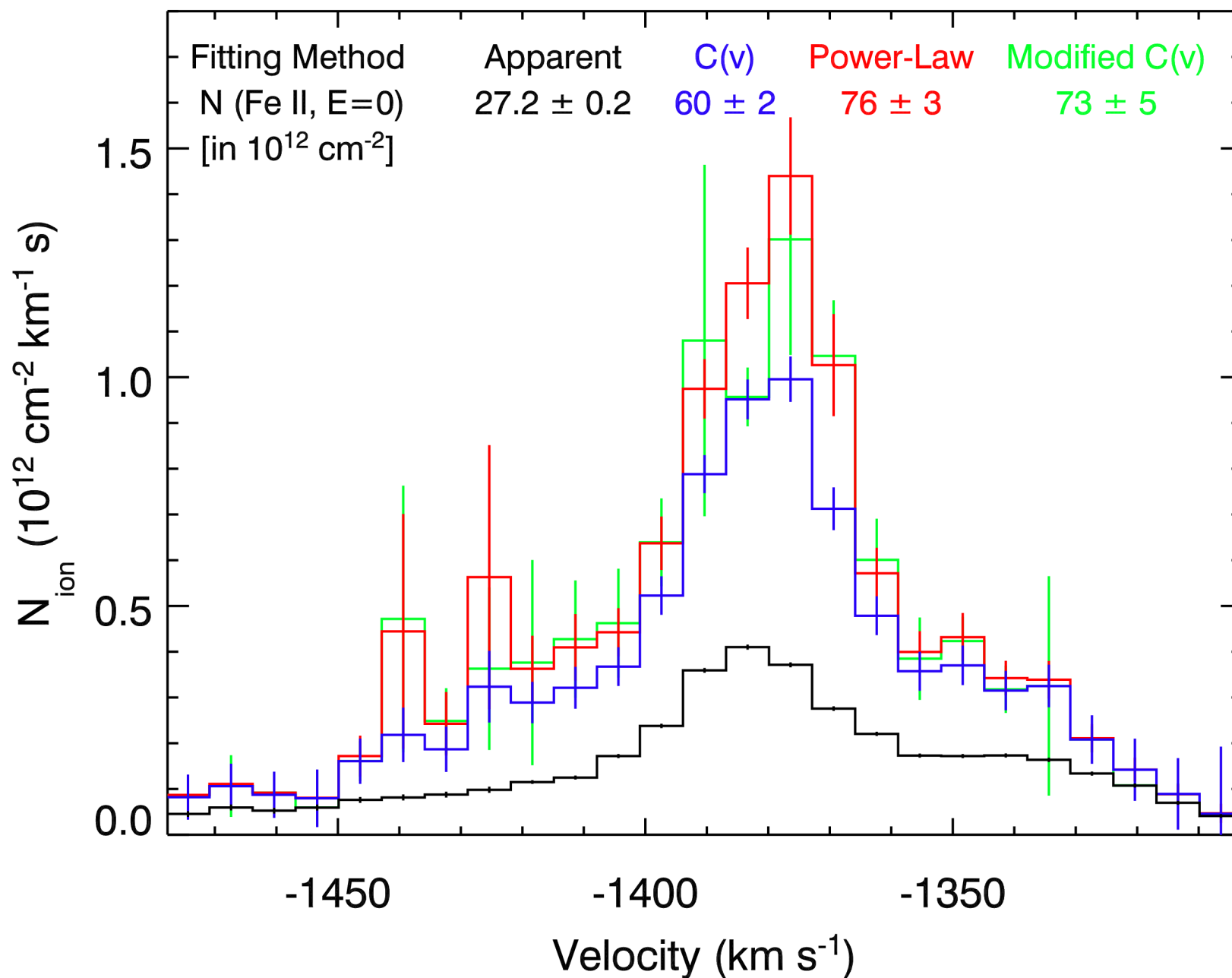
# Four ways to column densities



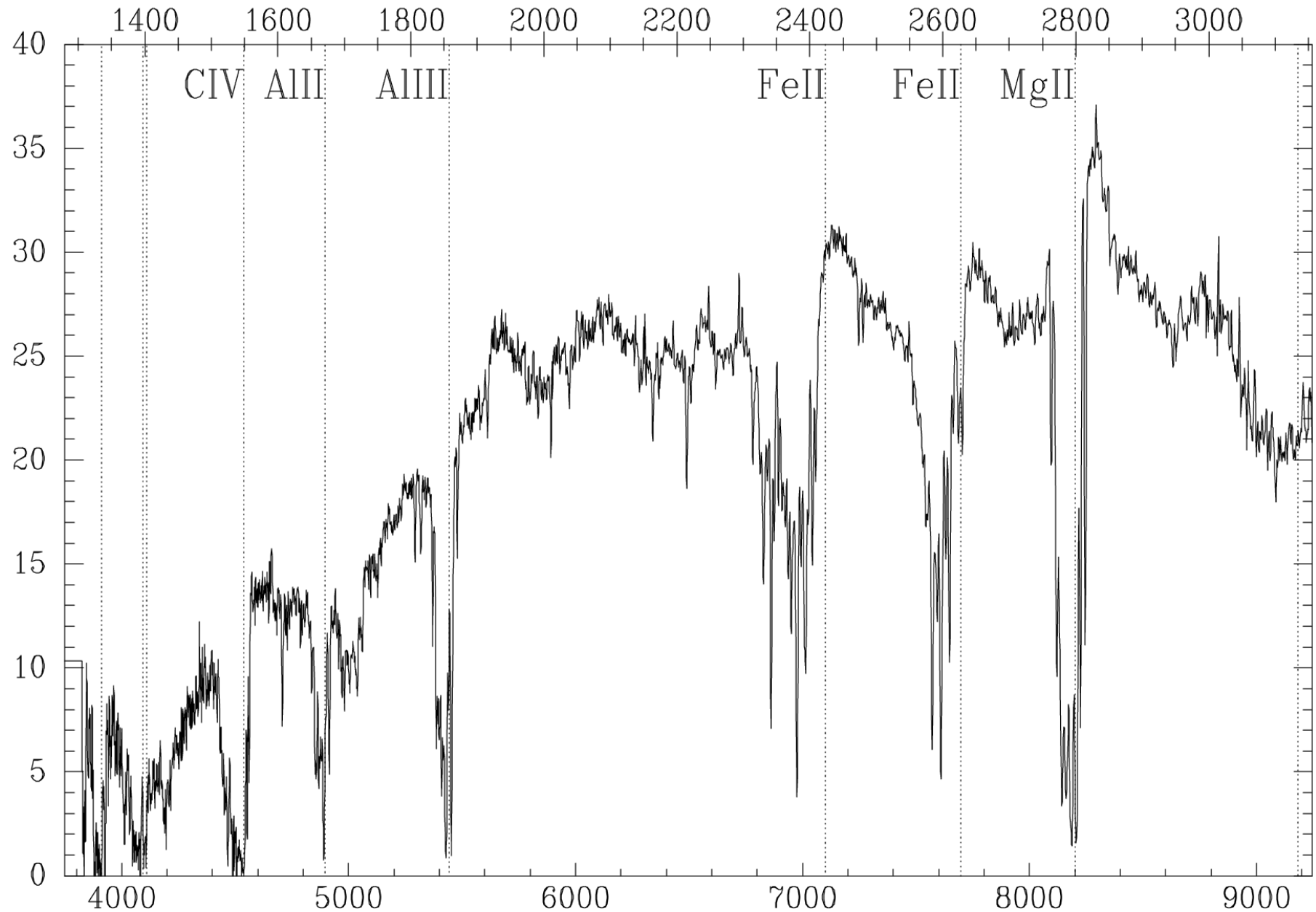
Fitting 3 Fe II lines 4 ways. From top: homogeneous, partial covering, power-law [best fit], modified partial covering.



# Resulting Fe II column as f(velocity)

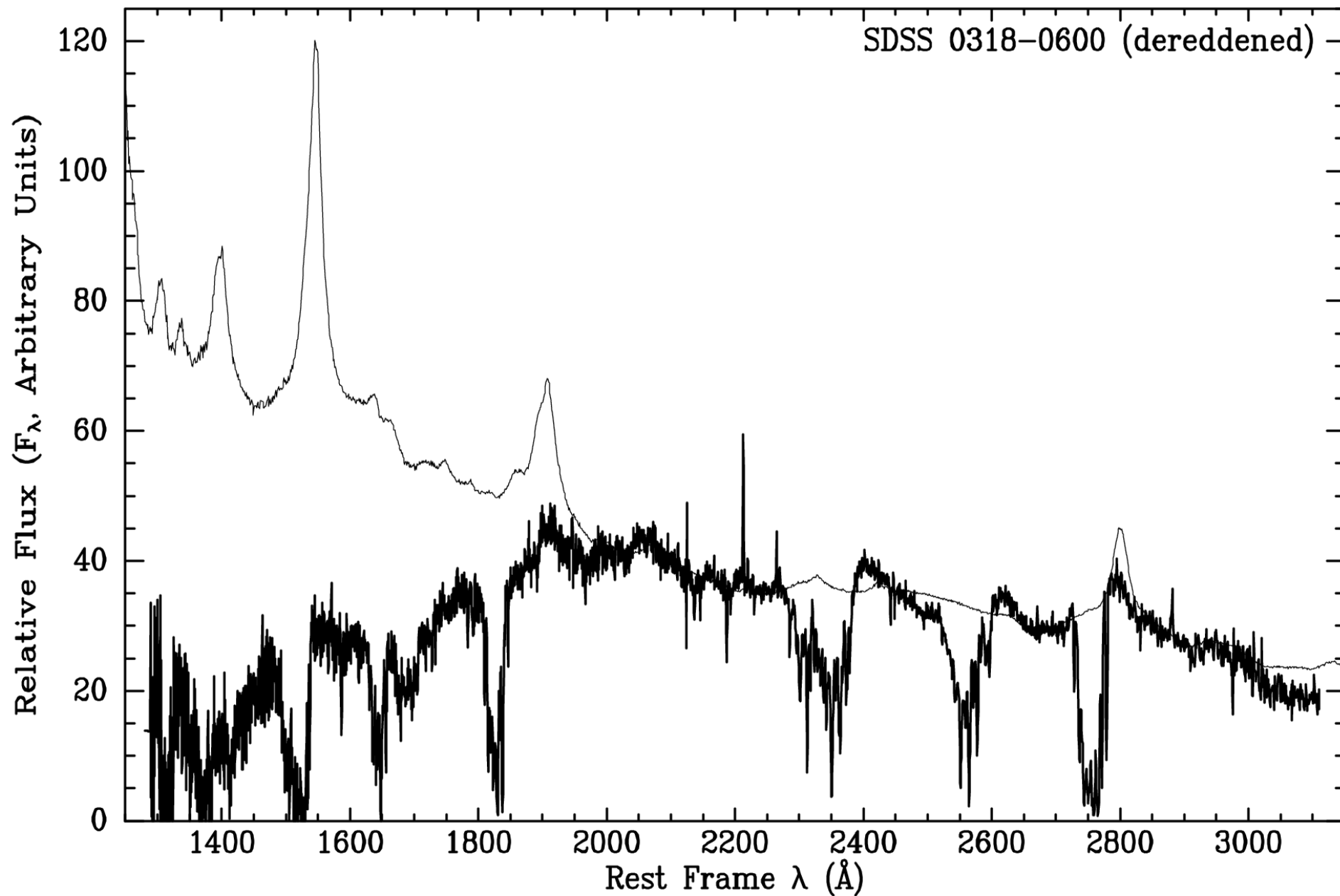


# SDSS spectrum of more complex BAL quasar (Dunn et al.)

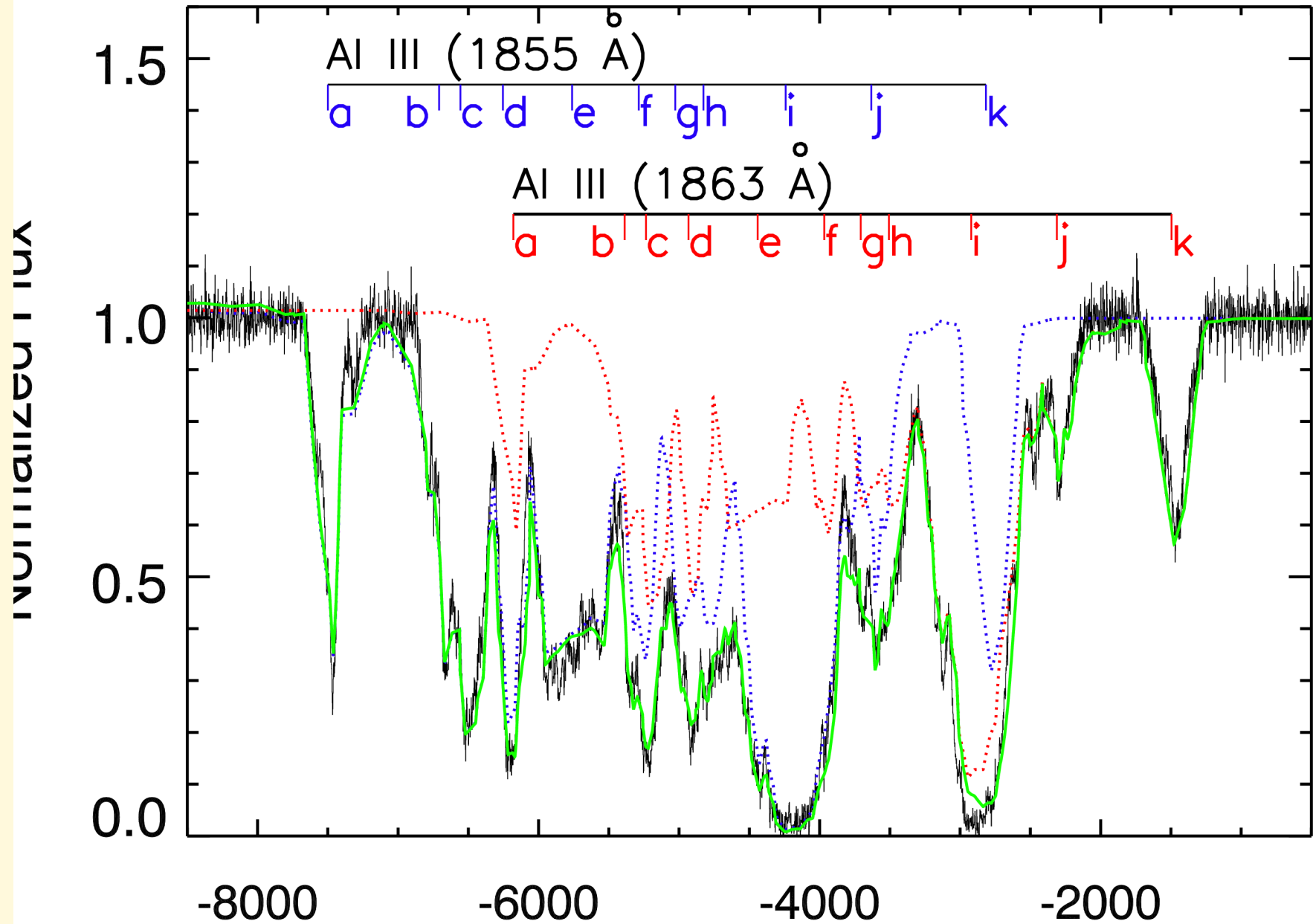


SDSSp J0318-0600  $z=1.93$  -- Bottom: Observed, Top: Rest

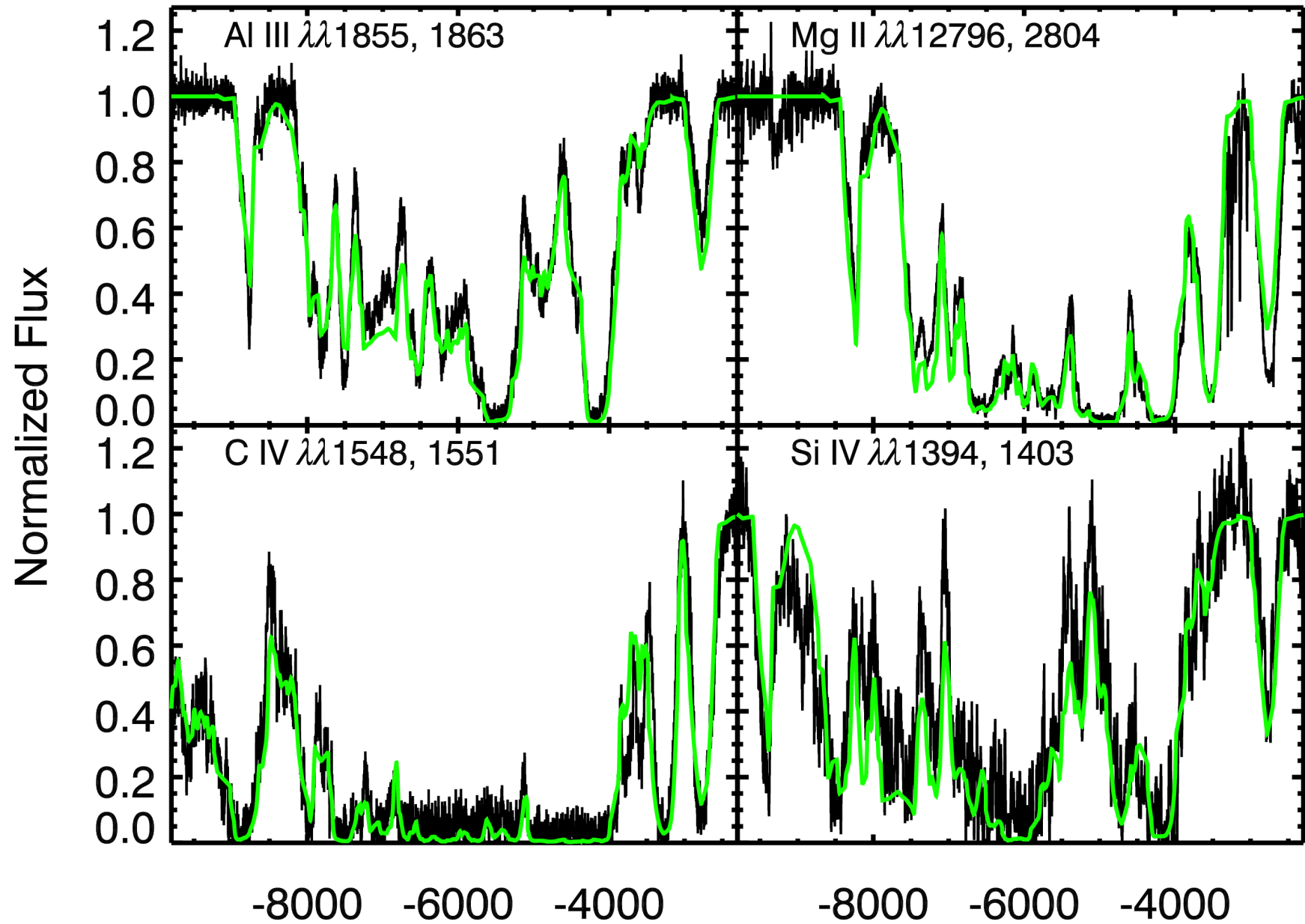
# Dereddened by SMC extinction curve



Use Al II singlet as templates for Al III doublet...



...and other lines (green fits to black data)





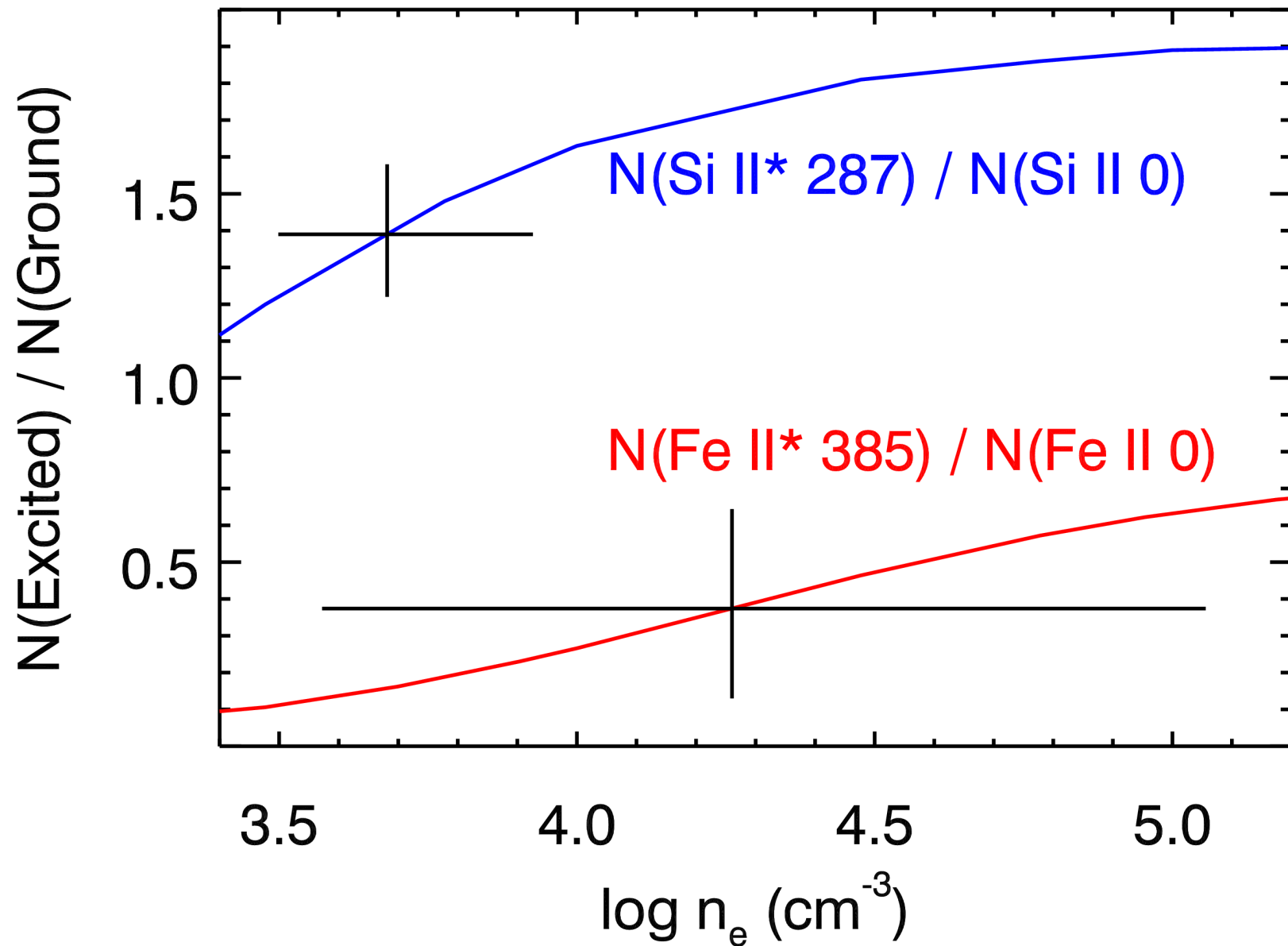
## From $N_{ion}$ to $N_H$

- Column density measurements  $N_{ion}$  are reasonably secure.
- Uncertainties: oscillator strengths, continuum placement, coverage of accretion disk vs. broad emission line region.
- Relate  $N_{ion}$  to  $N_H$  through photoionization modeling, for which a range of SEDs must be considered, and the hydrogen particle density  $n_H$  is needed as input.

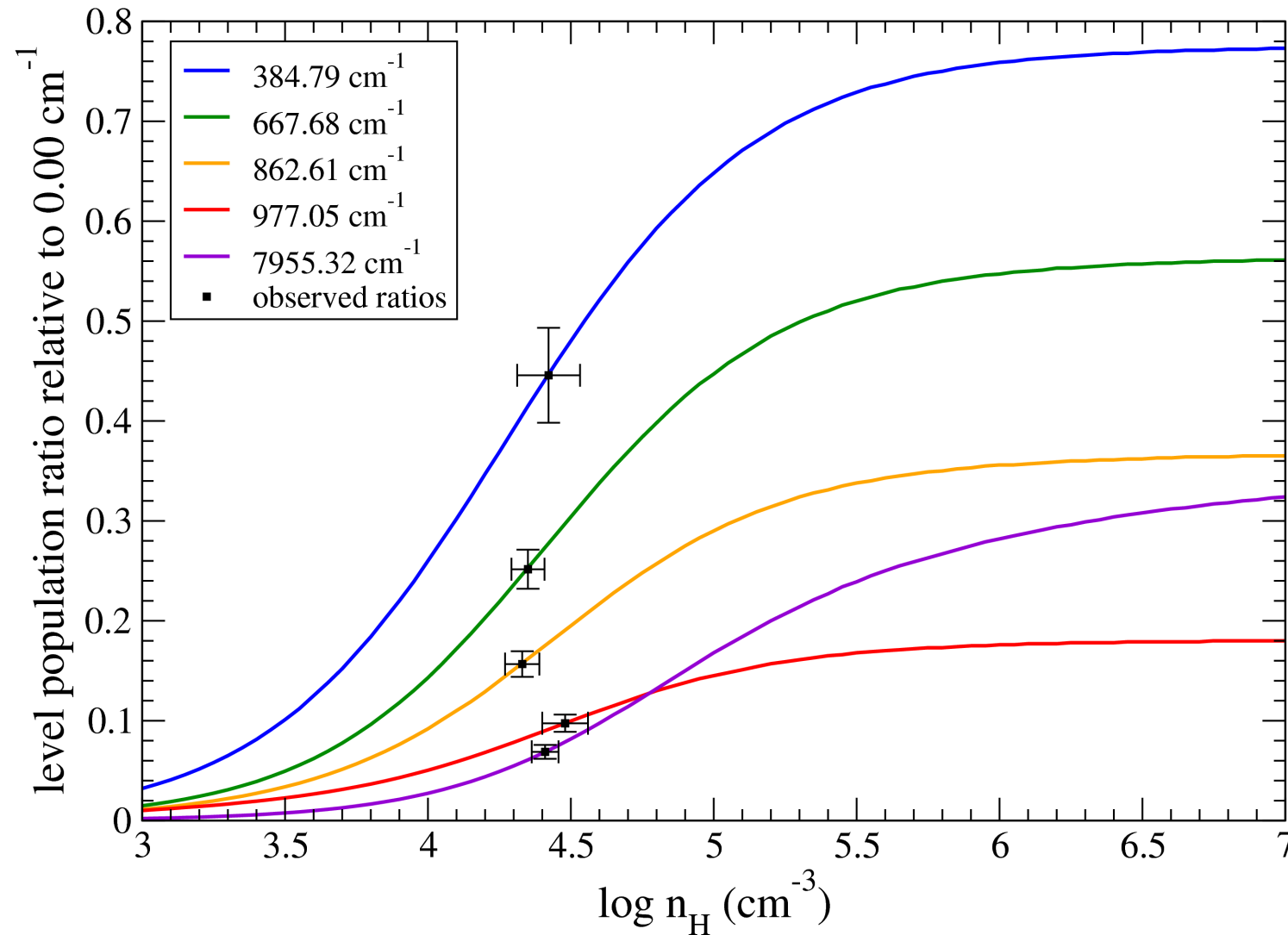
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- **Density constrained by looking for absorption from low-lying, metastable excited states (e.g., Fe II\*)**
- **The column density ratio of Fe II\* to Fe II increases rapidly near the critical  $n_e$  for that Fe II\* transition.**
- **Secondary dependences on temperature, radiative effects.**

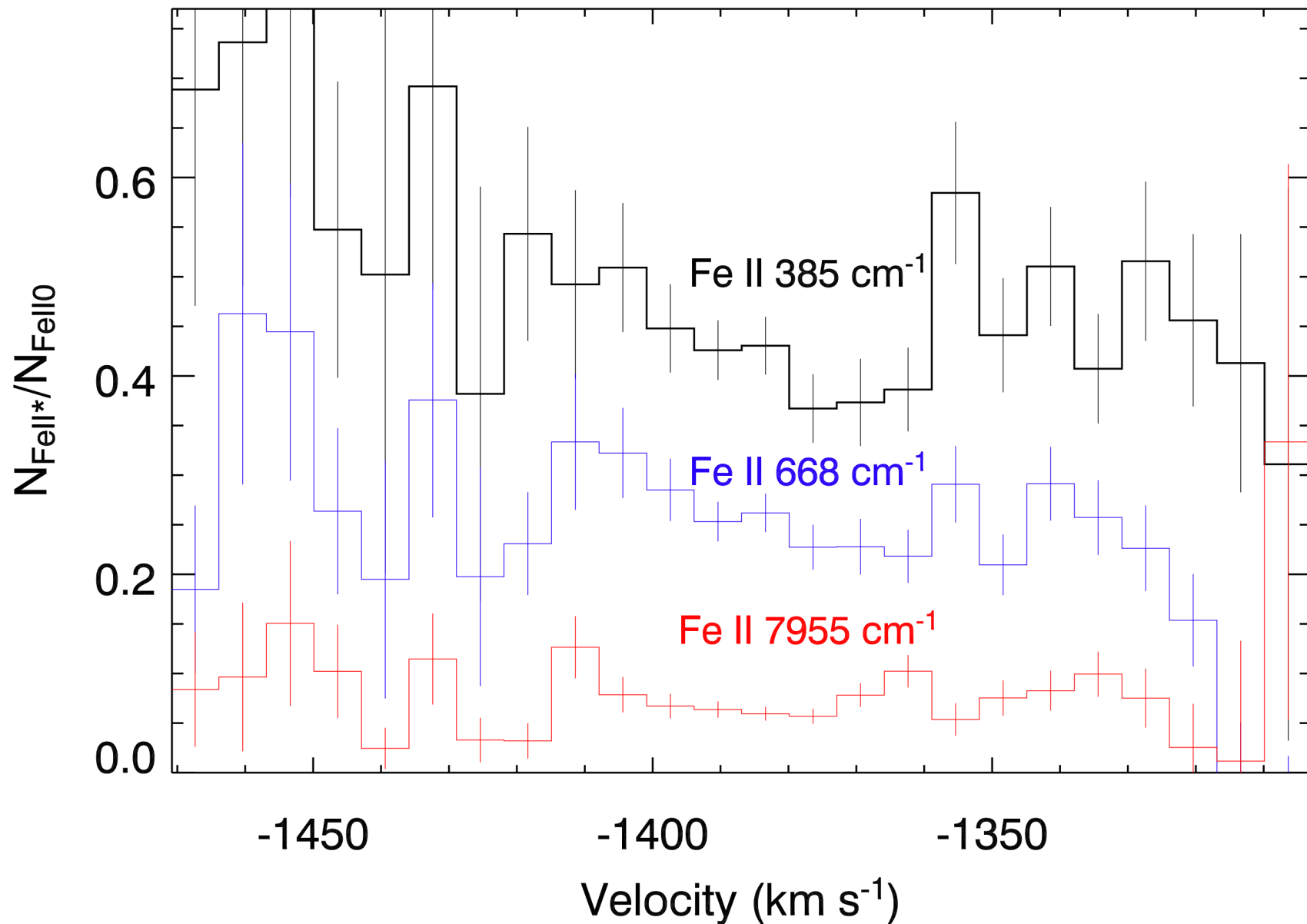
Density example:  $\log n_e = 3.75 \pm 0.22$  (Moe et al.)



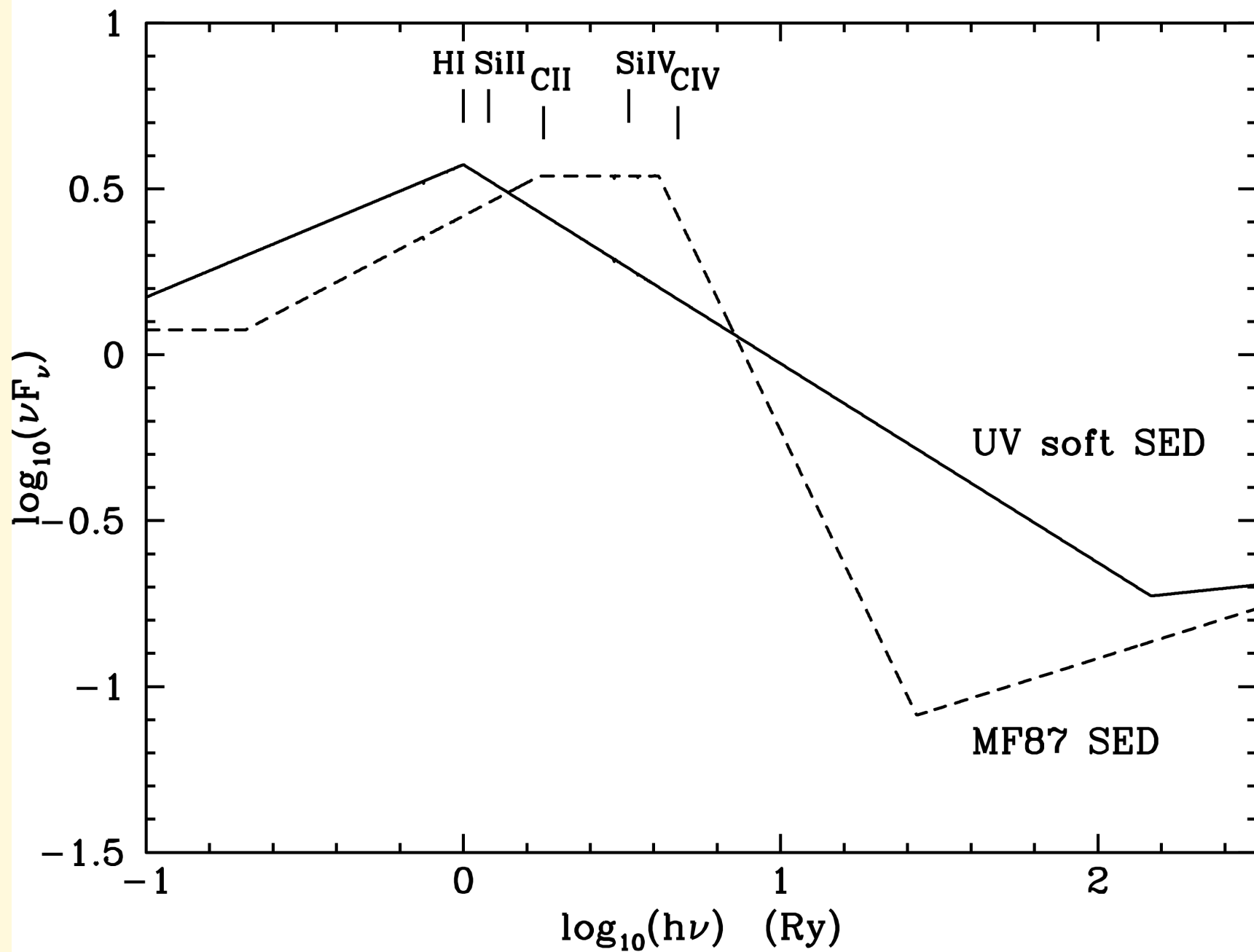
# Density example: $\log n_e = 4.4 \pm 0.1$ (Korista et al.)



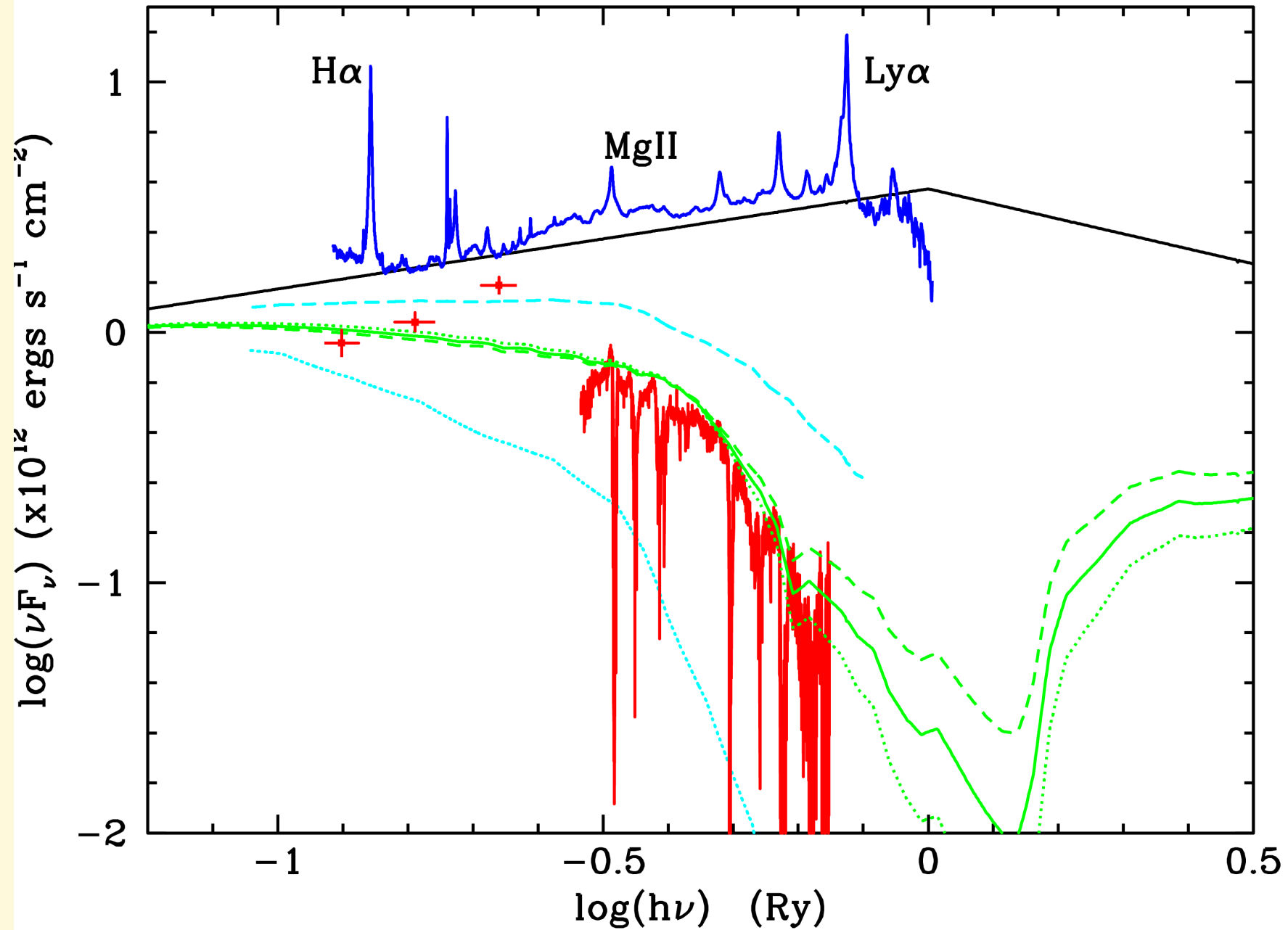
Density roughly constant with  $v$ , so sum  $N_{ion}$  over  $v$



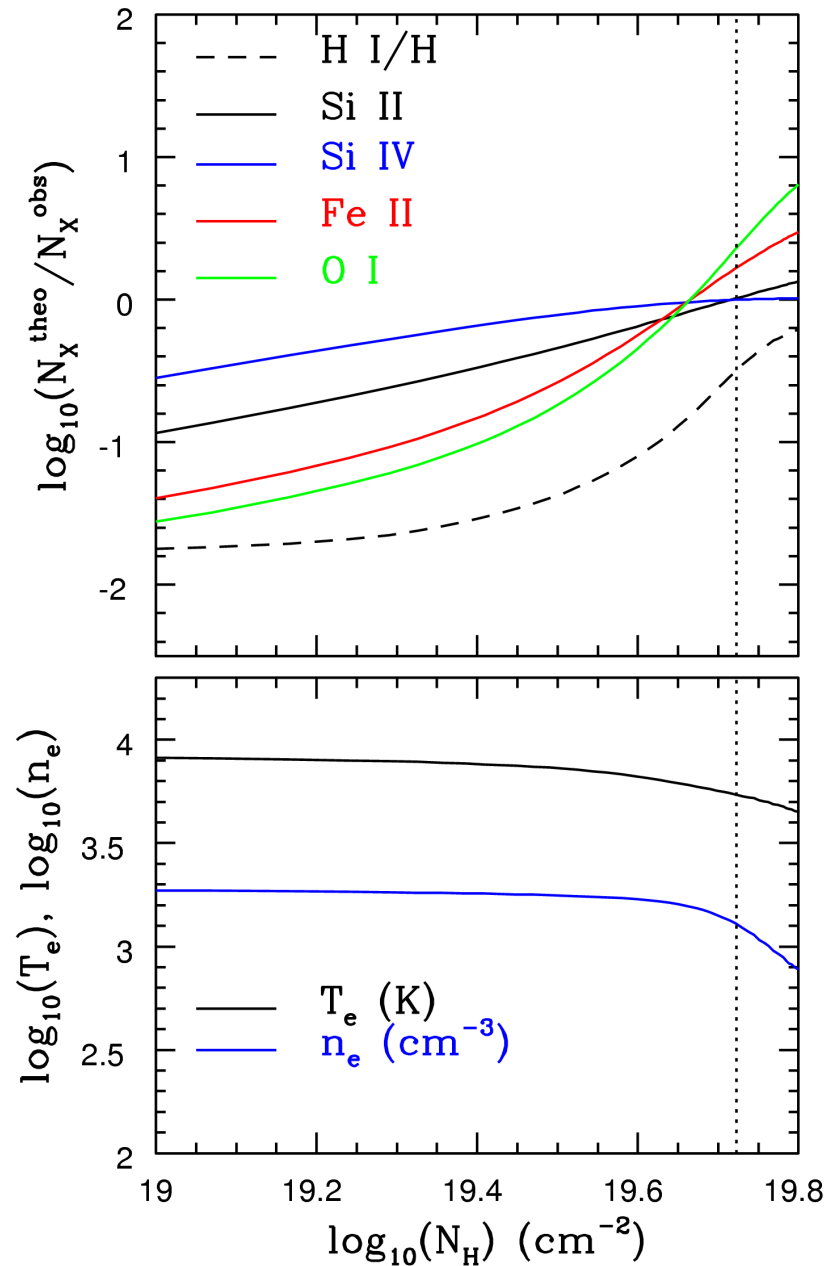
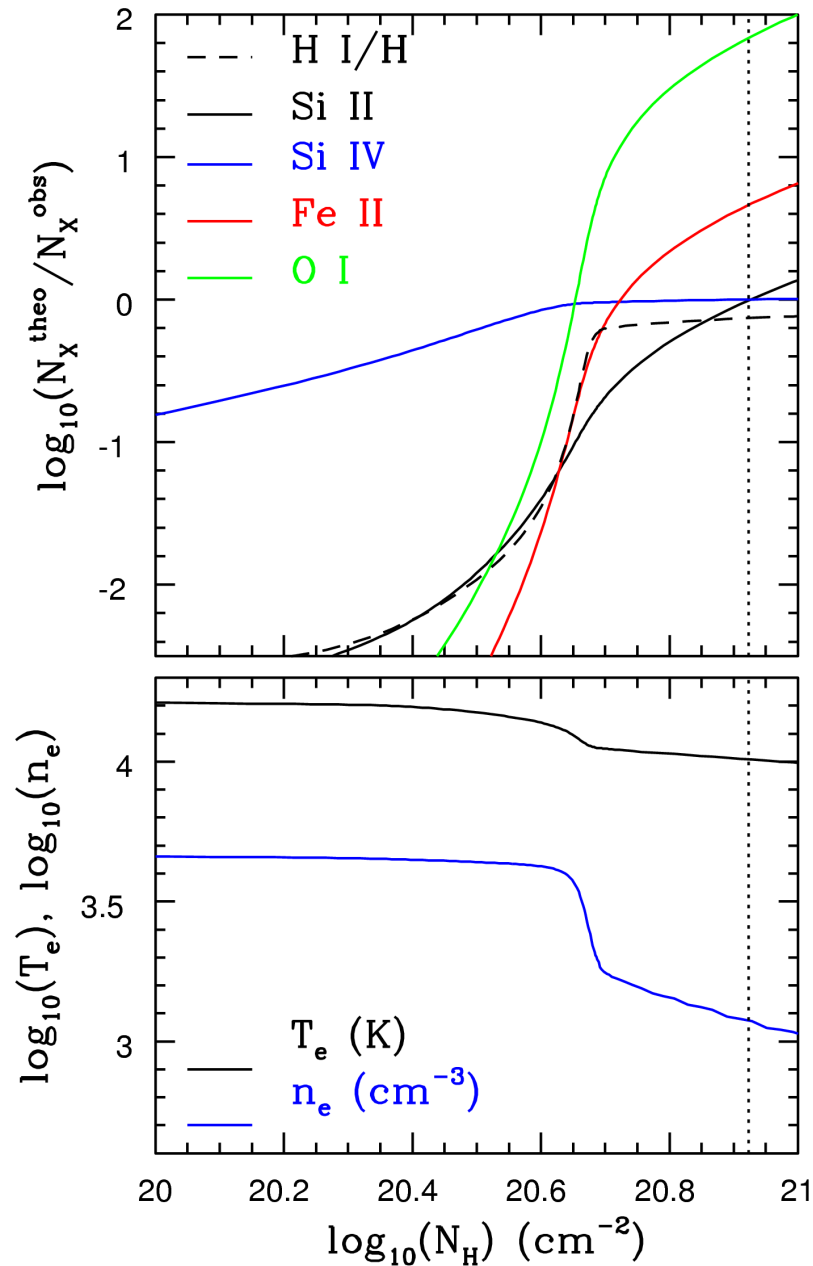
# SEDs studied in Dunn et al.



Attenuated & unattenuated SEDs also considered.



# Top panels: $N_{model}/N_{data}$ for $Z_{\odot}$ (left) & $7.2Z_{\odot}$ (right) models





## Results for Q0318-0600 (Dunn et al.)

- Best fit is attenuated SED,  
 $4.2 Z_{\odot}$ ,  $\log U = -3.02$ ,  $\log N_H = 20.1 \text{ cm}^{-2}$ ,  $R = 5.5 \text{ kpc}$
- Also acceptable fit from unattenuated SED,  
 $7.2 Z_{\odot}$ ,  $\log U = -2.85$ ,  $\log N_H = 19.9 \text{ cm}^{-2}$ ,  $R = 18.7 \text{ kpc}$
- $\dot{M} = 160\text{--}330 M_{\odot} \text{ yr}^{-1}$  ( $\Omega = 0.2$ )
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- $L_k/L_{bol} = 0.2\text{--}0.4\%$  ( $\Omega = 0.2$ )
- Other outflows have up to 10x higher  $L_k/L_{bol}$ , but still that's at most a few % of  $L_{bol}$  in kinetic luminosity.
- However, most studies to date done at low  $v < 5000 \text{ km/s}$
- Plus, any hotter phase of the outflow isn't sampled, and in Seyferts that can be a multiplier of 4–100 (Gabel et al 2005, Arav et al 2007)

Table 9. Properties of Measured Outflows to Date

Object	$R^a$ (kpc)	$\log N_H$ ( $\text{cm}^{-2}$ )	$\log U_H$	$\log \dot{E}_k$ ( $\text{ergs s}^{-1}$ )	$\dot{M}$ ( $M_\odot \text{ yr}^{-1}$ )	Reference <sup>b</sup>
QSO 0059-2735	0.001 - 0.05	$\gtrsim 21.5^c$	-0.7	$\gtrsim 41.1 - 42.8$	$\gtrsim 0.2$	1
3C 191	28	20.3	-2.8	44.0	310	2
QSO 1044+3656	0.1 - 2.1	20.0 - 22.0	-1.0 - -6.0	44.5 - 45.4	74 - 530	3
FIRST 1214+2803	0.001 - 0.03	21.4 - 22.2	-2.0 - -0.7	41.6 - 43.8	0.3 - 55	4
FIRST 0840+3633	0.001	$\sim 21.3$	$< -1.8$	$> 41.9$	$> 0.3$	5
FIRST 0840+3633 <sup>d</sup>	0.23	—	—	—	—	5
QSO 2359-1241	3	20.6	-2.4	43.7	93	6
SDSS J0838+2955	3.3	20.8	-1.9	45.7	590	7
SDSS J0318-0600	6 or 17	19.9 or 20.0	-3.1 or -2.7	44.8 or 45.4	120 or 450	8

<sup>a</sup>For relative accuracies, see Section 1.

<sup>b</sup>1-Wampler et al. (1995), 2-Hamann et al. (2001), 3-de Kool et al. (2001), 4-de Kool et al. (2002a), 5-de Kool et al. (2002b), 6-Korista et al. (2008), 7-Moe et al. (2009), 8-This Work

<sup>c</sup>Based on Table 5 in Wampler et al. (1995)

<sup>d</sup>Distance derived from Fe II fluorescence and no photoionization modeling was performed for this object

# Uncertainties in $R$

- $L_k = \frac{1}{2}\dot{M}v^2$  where  $\dot{M} \geq 4\pi\mu m_p v N_H R \Omega$
- C II\* & Si II\* have low critical densities
- Recall  $U_H = Q_H/4\pi R^2 c n_H$ : take an outflow of the observed low  $n_H$  and move it closer to quasar. As  $R$  decreases,  $U_H$  increases; eventually, low-ionization gas will disappear
- But at smaller  $R$ , higher-density gas can still have  $U_H$  low enough for Fe II to exist, and some Fe II\* lines have high critical densities
- Such higher-density tracers (incl. C III\*, Fe III\*) should probe to smaller distances; outflows at many scales?
- Separate issue: X-ray absorption can modify spectrum, reduce distance (Everett et al 2002)

# Uncertainties in $\Omega$

- $L_k = \frac{1}{2}\dot{M}v^2$  where  $\dot{M} \geq 4\pi\mu m_p v N_H R \Omega$
- What  $\Omega$  to use?  $\Omega_{obs} = 0.2$ : 20% of our sightlines to quasars have BALs in them, but if 50% of quasars are obscured, BALs cover only 10% of the sky as seen from the quasar ( $\Omega \leq f_{BAL}$ )
- To date, distance measurements made only for the 1 in 100 BAL quasars with Fe II\*, C II\*, Si II\*. So,  $\Omega = 0.002$ ?
- No. But fair to ask how similar are excited-state sightlines to more typical sightlines (answer: a few times higher column).
- Test by looking at S IV\* / S IV (Dunn in prep.); same ionization as C IV, detected at rate consistent with same  $\Omega = 0.2$
- Regardless, need many objects to get average outflow picture

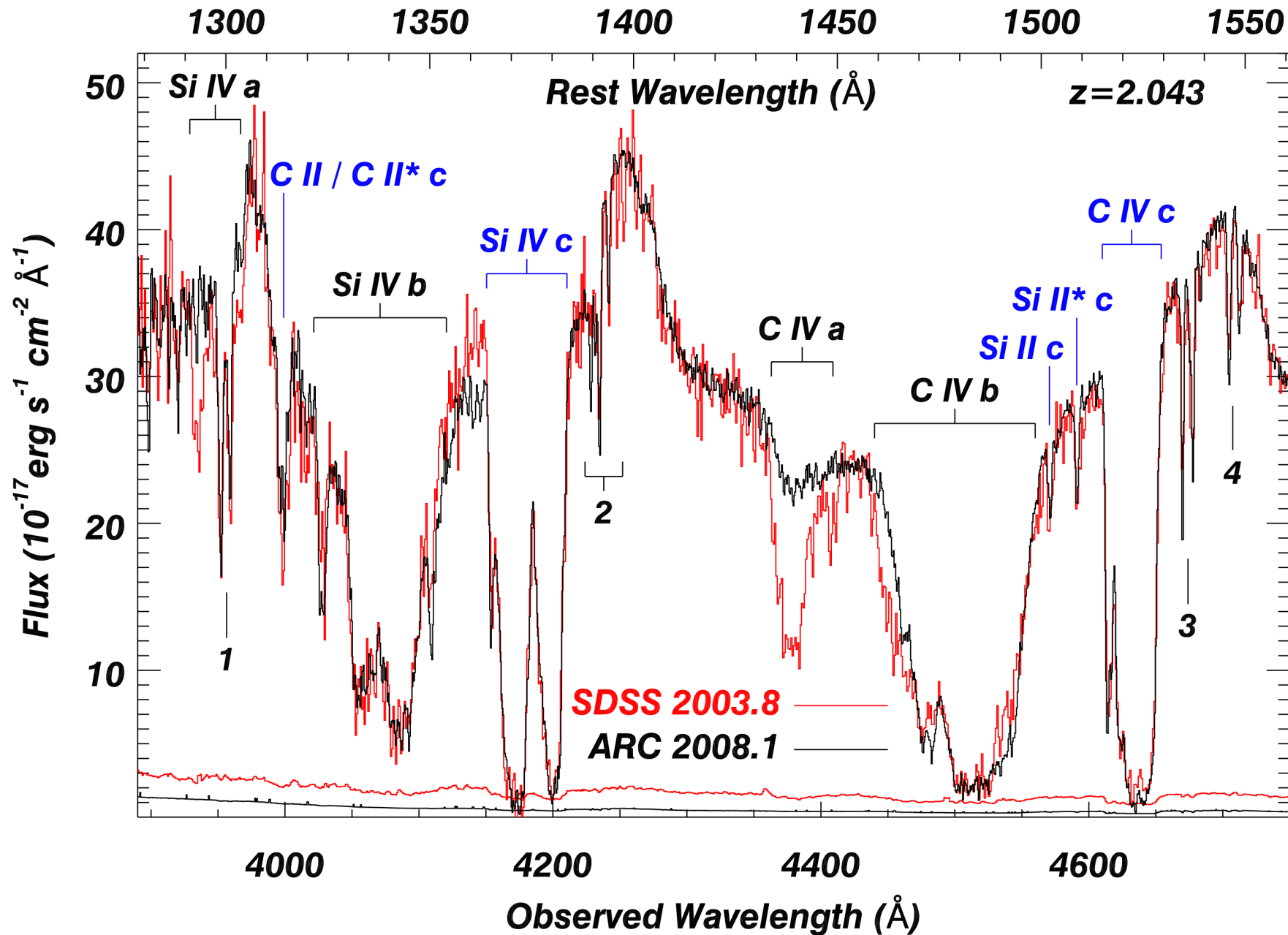
# Conclusions

- Ionic column densities can be measured if care is taken (partial covering or more sophisticated models)
- To date, C II & Si II used to probe low  $n_e$  and thus preferentially larger distances, but that is changing
- Photoionization modeling yields  $N_H$ , but I would like to see a wider range of models explored (e.g.: continuous wind; physical model for location of X-ray obscuration)
- Nevertheless, some outflows are tens of kpc away from the BH that launched them (3C 191)
- Some are only few pc away (Hall et al. arXiv next week)
- Atomic data often a limiting factor

# QUASARS



# Moe et al. (2009) quasar SDSS J0838+2955





# Moe et al. (2009) joint $U_H$ , $N_H$ constraints

