

IV. Laplace transform

Given a function $f(t)$ defined for $0 \leq t < \infty$, the Laplace transform is

$$\mathcal{L}[f(t)] = F(s) = \int_0^{\infty} dt e^{-st} f(t)$$

This is an example of an integral transform.

$F(s)$ is only defined as long as the integral converges.

ex. $f(t) = c = \text{constant}$

$$F(s) = \int_0^{\infty} dt c e^{-st} = -\frac{c}{s} e^{-st} \Big|_0^{\infty} = \frac{c}{s}$$

ex. $f(t) = t$

$$F(s) = \int_0^{\infty} dt t e^{-st} = -\frac{1}{s} t e^{-st} \Big|_0^{\infty} + \int_0^{\infty} dt \frac{1}{s} e^{-st} = \frac{1}{s^2}$$

ex. $f(t) = t^n$ ($n=0, 1, 2, \dots$)

$$\begin{aligned} F(s) &= \int_0^{\infty} dt t^n e^{-st} = -\frac{1}{s} t^n e^{-st} \Big|_0^{\infty} + \int_0^{\infty} dt n t^{n-1} \frac{1}{s} e^{-st} \\ &= n(n-1)(n-2)\dots(1) \frac{1}{s^n} \int_0^{\infty} dt e^{-st} = \frac{n!}{s^{n+1}} \end{aligned}$$

converges for $s \geq 0$, so $F(s)$ defined for $s \geq 0$.

ex. $f(t) = e^{\alpha t}$

$$F(s) = \int_0^{\infty} dt e^{(\alpha-s)t} = \frac{1}{\alpha-s} e^{(\alpha-s)t} \Big|_0^{\infty} = \frac{1}{s-\alpha}$$

but only converges for $s \geq \alpha$. So $F(s)$ defined for $s \geq \alpha$.

$$\text{ex. } f(t) = \cos(\omega t) = \frac{1}{2} (e^{i\omega t} + e^{-i\omega t})$$

$$\begin{aligned} F(s) &= \int_0^{\infty} dt \cos(\omega t) e^{-st} \\ &= \frac{1}{2} \int_0^{\infty} dt (e^{(i\omega-s)t} + e^{(-i\omega-s)t}) \\ &= \frac{1}{2} \left[\frac{1}{i\omega-s} e^{(i\omega-s)t} + \frac{1}{-i\omega-s} e^{(-i\omega-s)t} \right] \Big|_{t=0}^{\infty} \\ &= \frac{1}{2} \left(\frac{1}{s-i\omega} + \frac{1}{s+i\omega} \right) = \frac{s+i\omega+s-i\omega}{2(s^2+\omega^2)} = \frac{s}{s^2+\omega^2} \end{aligned}$$

$$\text{ex. } f(t) = \sin(\omega t)$$

$$\text{F(s)} = \frac{\omega}{s^2 + \omega^2}$$

$F(s)$ for $\sin(\omega t)$ & $\cos(\omega t)$ defined for $s \geq 0$.

The Laplace transform is a linear operator:

$$(1) \mathcal{L}[c f(t)] = \int_0^{\infty} dt e^{-st} c f(t) = c \int_0^{\infty} dt e^{-st} f(t) = c \mathcal{L}[f(t)]$$

$$\begin{aligned} (2) \mathcal{L}[f_1(t) + f_2(t)] &= \int_0^{\infty} dt e^{-st} (f_1(t) + f_2(t)) = \int_0^{\infty} dt e^{-st} f_1(t) + \int_0^{\infty} dt e^{-st} f_2(t) \\ &= \mathcal{L}[f_1(t)] + \mathcal{L}[f_2(t)] \end{aligned}$$

Laplace transform of derivatives:

$$\begin{aligned} \mathcal{L}[f'(t)] &= \int_0^{\infty} dt e^{-st} \frac{df}{dt} = \left[-\frac{1}{s} e^{-st} f(t) \right]_0^{\infty} + s \int_0^{\infty} dt e^{-st} f(t) \\ &= -f(0) + s F(s) \end{aligned}$$

$$\begin{aligned}\mathcal{L}[f''(t)] &= -f'(0) + s\mathcal{L}[f'(t)] \\ &= -f'(0) - sf(0) + s^2 F(s)\end{aligned}$$

So an arbitrary number of derivatives gives:

$$\mathcal{L}[f^{(n)}(t)] = s^n F(s) - f^{(n-1)}(0) - sf^{(n-2)}(0) - \dots - s^{n-1} f(0)$$

Using Laplace transforms to solve differential equations:

Useful for systems that begin with an initial condition at time $t=0$ ~~and then~~ (or some other time) and then evolve for $t>0$.

ex. simple harmonic oscillator $\ddot{x}(t) + \omega^2 x(t) = 0$

$$\text{define } X(s) = \mathcal{L}[x(t)] = \int_0^\infty dt e^{-st} x(t)$$

Take Laplace transform of differential equation:

$$\mathcal{L}[\ddot{x}(t) + \omega^2 x(t)] = s^2 X(s) - x'(0) - sx(0) + \omega^2 X(s) = 0$$

$$\begin{aligned}\Rightarrow X(s) &= \frac{x'(0) + sx(0)}{s^2 + \omega^2} \\ &= x'(0) \frac{1}{s^2 + \omega^2} + x(0) \frac{s}{s^2 + \omega^2}\end{aligned}$$

$$x(t) = \frac{x'(0)}{\omega} \sin(\omega t) + x(0) \cos(\omega t)$$

Laplace transform transforms a differential equation into an algebraic equation.

example: radioactive decay

Consider a chain of radioactive decays $1 \rightarrow 2 \rightarrow 3$, where $1, 2, 3$ are elements and $\lambda_1, \lambda_2, \lambda_3$ are the decay rates. ($\lambda_3 = 0 \rightarrow 3$ is stable.)

Let n_1, n_2, n_3 be the number of a given element.
Coupled equations:

$$\frac{dn_1}{dt} = -\lambda_1 n_1 \quad \frac{dn_2}{dt} = -\lambda_2 n_2 + \lambda_1 n_1 \quad \frac{dn_3}{dt} = +\lambda_2 n_2.$$

take Laplace transform:

~~$$sN_1(s) - n_1(0) = -\lambda_1 N_1(s)$$~~

$$sN_1(s) - n_1(0) = -\lambda_1 N_1(s)$$

$$sN_2(s) - n_2(0) = -\lambda_2 N_2(s) + \lambda_1 N_1(s)$$

$$sN_3(s) - n_3(0) = \lambda_2 N_2(s)$$

Solve algebraic equations:

$$N_1(s) = \frac{n_1(0)}{s + \lambda_1}$$

$$N_2(s) = \frac{n_2(0)}{s + \lambda_2} + \frac{\lambda_1 N_1(s)}{s + \lambda_2} = \frac{n_2(0)}{s + \lambda_2} + \frac{\lambda_1 n_1(0)}{(s + \lambda_1)(s + \lambda_2)}$$

$$N_3(s) = \frac{n_3(0)}{s} + \frac{\lambda_2 N_2(s)}{s} \quad \text{~~to } N_2(s)~~$$

$$= \frac{n_3(0)}{s} + \frac{\lambda_2}{s} \left(\frac{n_2(0)}{s + \lambda_2} + \frac{\lambda_1 n_1(0)}{(s + \lambda_1)(s + \lambda_2)} \right)$$

$$N_2(s) = \frac{n_2(0)}{s+\lambda_2} + \frac{\lambda_1 n_1(0)}{\lambda_2 - \lambda_1} \left(\frac{1}{s+\lambda_1} - \frac{1}{s+\lambda_2} \right)$$

$$= \left(n_2(0) - \frac{\lambda_1 n_1(0)}{\lambda_2 - \lambda_1} \right) \frac{1}{s+\lambda_2} + \frac{\lambda_1 n_1(0)}{\lambda_2 - \lambda_1} \frac{1}{s+\lambda_1}$$

~~$$N_2(s) = \frac{n_2(0)}{s} + \frac{n_2(0)}{s} - \frac{n_2(0)}{s+\lambda_2}$$~~

$$N_3(s) = \frac{n_3(0)}{s} + n_1(0) \left(\frac{1}{s} + \frac{\lambda_2}{\lambda_1 - \lambda_2} \frac{1}{s+\lambda_1} - \frac{\lambda_1}{\lambda_1 - \lambda_2} \frac{1}{s+\lambda_2} \right)$$

$$= \frac{1}{s} (n_1(0) + n_2(0) + n_3(0)) + n_1(0) \frac{\lambda_2}{\lambda_1 - \lambda_2} \frac{1}{s+\lambda_1}$$

$$- \left(n_2(0) + \frac{n_1(0) \lambda_1}{\lambda_1 - \lambda_2} \right) \frac{1}{s+\lambda_2}$$

Next, invert the transform:

So $n_1(t) = n_1(0) e^{-\lambda_1 t}$

$$n_2(t) = \left(n_2(0) + \frac{\lambda_1 n_1(0)}{\lambda_1 - \lambda_2} \right) e^{-\lambda_2 t} + \frac{\lambda_1 n_1(0)}{\lambda_2 - \lambda_1} e^{-\lambda_1 t}$$

$$n_3(t) = n_1(0) + n_2(0) + n_3(0) + n_1(0) \frac{\lambda_2}{\lambda_1 - \lambda_2} e^{-\lambda_1 t}$$

$$- \left(n_2(0) + \frac{n_1(0) \lambda_1}{\lambda_1 - \lambda_2} \right) e^{-\lambda_2 t}$$

$$n_2(t) = n_2(0) e^{-\lambda_2 t} + n_1(0) \frac{\lambda_1}{\lambda_2 - \lambda_1} (e^{-\lambda_1 t} - e^{-\lambda_2 t})$$

$$n_3(t) = n_3(0) + n_2(0) (1 - e^{-\lambda_2 t})$$

$$+ n_1(0) \left(1 + \frac{\lambda_2}{\lambda_1 - \lambda_2} e^{-\lambda_1 t} - \frac{\lambda_1}{\lambda_1 - \lambda_2} e^{-\lambda_2 t} \right)$$

Inverse Laplace transform

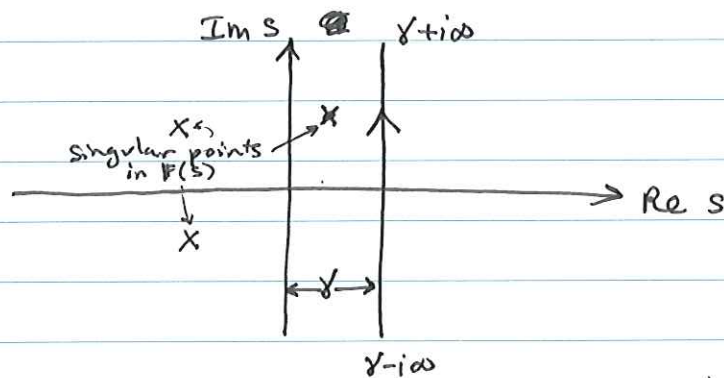
Easy way to invert the Laplace transform is by inspection
(i.e. reducing $F(s)$ to a form where you can guess $f(t)$)

More general procedure: Bromwich integral

Given $F(s)$, the inverse Laplace transformed function $f(t)$ is

$$f(t) = \frac{1}{2\pi i} \int_{\gamma-i\infty}^{\gamma+i\infty} ds e^{st} F(s)$$

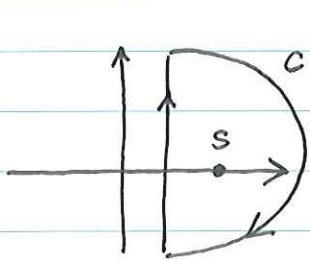
where γ is any real number chosen so that the contour lies to the right of any ~~poles~~ singularities in $F(s)$.



proof: take the Laplace transform

$$\begin{aligned} \int_0^{\infty} dt f(t) e^{-st} &= \int_0^{\infty} dt e^{-st} \frac{1}{2\pi i} \int_{\gamma-i\infty}^{\gamma+i\infty} ds' e^{s't} F(s') \\ &= \frac{1}{2\pi i} \int_{\gamma-i\infty}^{\gamma+i\infty} ds' F(s') \int_0^{\infty} dt e^{-(s-s')t} \quad \uparrow \text{if } s > \gamma, \text{ then integral converges} \\ &= \frac{1}{2\pi i} \int_{\gamma-i\infty}^{\gamma+i\infty} ds' \frac{-F(s')}{s' - s} \end{aligned}$$

Now close the contour in $\text{Re}(s) > 0$ half-plane:



$$\int_{\gamma-i\infty}^{\gamma+i\infty} ds'$$

$$= \oint_C ds'$$

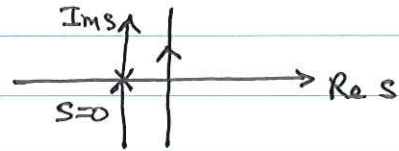
provided $F(s') \rightarrow 0$ for $s' \rightarrow \infty$.

$$\Rightarrow \frac{1}{2\pi i} \int_{\gamma-i\infty}^{\gamma+i\infty} ds' \frac{F(s')}{s-s'} = \frac{1}{2\pi i} \oint_C ds' \frac{-F(s')}{s'-s}$$

$$= \frac{1}{2\pi i} (-2\pi i) (-F(s)) = F(s)$$

example: $F(s) = \frac{1}{s^2}$. pole at $s=0 \rightarrow$ choose $\gamma > 0$.

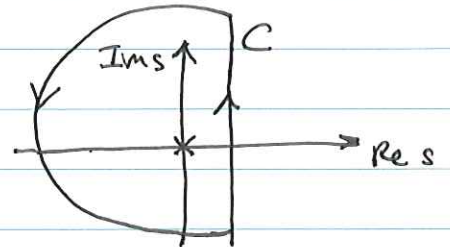
$$f(t) = \frac{1}{2\pi i} \int_{\gamma-i\infty}^{\gamma+i\infty} ds e^{st} \frac{1}{s^2}$$



Note: ~~the contour~~ closing the contour at $\text{Re}(s) \rightarrow \infty$ will not converge. Need to close the contour at $\text{Re}(s) \rightarrow -\infty$.

$$f(t) = \frac{1}{2\pi i} \oint_C ds e^{st} \frac{1}{s^2}$$

where



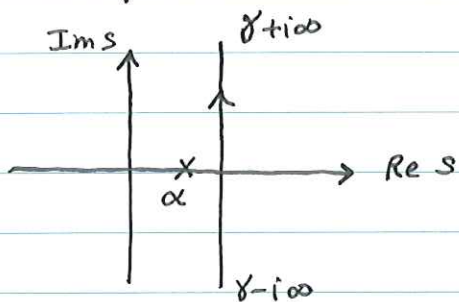
~~Residue~~

$$f(t) = \frac{1}{2\pi i} \times 2\pi i \text{Res} \left[e^{st} \frac{1}{s^2}; s=0 \right]$$

$$= \lim_{s \rightarrow 0} \frac{d}{ds} \left(s^2 \frac{e^{st}}{s^2} \right) = \lim_{s \rightarrow 0} t e^{st} = t$$

so $f(t) = t$.

example: $F(s) = \frac{1}{s-\alpha}$ (where $\alpha > 0$ is a real number)



Must choose $\gamma > \alpha$.

$$f(t) = \frac{1}{2\pi i} \int_{\gamma-i\infty}^{\gamma+i\infty} ds e^{st} \frac{1}{s-\alpha} = \frac{1}{2\pi i} \oint_C ds \frac{e^{st}}{s-\alpha}$$

close contour at $\text{Re}(s) \rightarrow -\infty$.

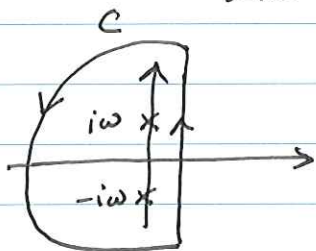
pole at $s = \alpha$.

$$f(t) = 2\pi i \frac{1}{2\pi i} \text{Res} \left[\frac{e^{st}}{s-\alpha}; s = \alpha \right] = e^{\alpha t}$$

example: $F(s) = \frac{s}{s^2 + \omega^2} = \frac{s}{(s+i\omega)(s-i\omega)}$ ($\omega > 0$, real)

choose $\gamma > \text{Re}(i\omega) = 0$

$$f(t) = \frac{1}{2\pi i} \int_{\gamma-i\infty}^{\gamma+i\infty} ds e^{st} \frac{s}{(s+i\omega)(s-i\omega)} = \frac{1}{2\pi i} \oint_C ds \frac{s e^{st}}{(s+i\omega)(s-i\omega)}$$



$$\begin{aligned} &= \text{Res} \left[\frac{s e^{st}}{(s+i\omega)(s-i\omega)}; s = i\omega \right] \\ &+ \text{Res} \left[\frac{s e^{st}}{s^2 + \omega^2}; s = -i\omega \right] \\ &= \frac{i\omega e^{i\omega t}}{2i\omega} + \frac{i\omega e^{-i\omega t}}{-2i\omega} = \cos(\omega t) \end{aligned}$$

Dirac delta function

Defined by $\delta(x-a) = \begin{cases} \infty & x=a \\ 0 & x \neq a \end{cases}$

and $\int_{-\infty}^{\infty} dx \delta(x-a) = 1.$

Some useful properties:

δ -function is even: $\delta(x-a) = \delta(a-x)$

$$\delta(c(x-a)) = \frac{1}{|c|} \delta(x-a)$$

$$\int_{-\infty}^{\infty} dx f(x) \delta(x-a) = f(a)$$

The δ -function is the derivative of the step function:

$$\theta(x-a) = \begin{cases} 1 & x > a \\ 0 & x < a \end{cases}$$

$$= \int_{-\infty}^x dx' \delta(x'-a)$$

$$\frac{d}{dx} \theta(x-a) = \delta(x-a)$$

Laplace transform:

let $f(t) = \delta(t-\tau)$

then $F(s) = \int_0^{\infty} dt e^{-st} \delta(t-\tau) = e^{-s\tau}.$

Green's functions

Very powerful technique for solving inhomogeneous, linear differential equations. Example: forced simple harmonic oscillator

$$\ddot{x} + 2\gamma \dot{x} + \omega^2 x = F(t)$$

linear = each term on the left side only involves power of $x(t)$ or its derivative

inhomogeneous = forcing function $F(t) \neq 0$.

~~Consider the case where the system is at rest for $t < 0$:~~

Consider the case where the system is at rest for $t < 0$:
i.e. $x(t) = \dot{x}(t) = F(t) = 0$ for $t < 0$.

Then ~~the~~ $F(t)$ "turns on" at $t = 0$. Solve for $x(t)$ with initial condition $x(0) = \dot{x}(0) = 0$.

Trick: first consider a simpler equation

$$\ddot{g} + 2\gamma \dot{g} + \omega^2 g = \delta(t - \tau) \quad (\text{where } \tau > 0)$$

The solution to this equation $g(t, \tau)$ is a Green's function.

After solving for $g(t, \tau)$, the solution for $x(t)$ is:

$$x(t) = \int_0^{\infty} d\tau g(t, \tau) F(\tau)$$

Proof: plug this solution into the original equation:

$$\begin{aligned} \left(\frac{d^2}{dt^2} + 2\gamma \frac{d}{dt} + \omega^2 \right) x(t) &= \int_0^{\infty} d\tau \left(\frac{d^2}{dt^2} + 2\gamma \frac{d}{dt} + \omega^2 \right) g(t, \tau) F(\tau) \\ &= \int_0^{\infty} d\tau \delta(t - \tau) F(\tau) = F(t) \end{aligned}$$

The Green's function describes the response of the system to an impulse at $t = \tau$.

The total response $x(t)$ to an arbitrary driving force $F(t)$ is the "sum" (integral) over all infinitesimal impulses for $0 < \tau < \infty$.

Now, solve for $g(t, \tau)$ using Laplace transform:

$$G(s, \tau) = \int_0^{\infty} dt e^{-st} g(t, \tau)$$

So we have, acting with $\int_0^{\infty} dt e^{-st}$ on both sides:

$$\begin{aligned} (s^2 G(s, \tau) - s g(0, \tau) - \dot{g}(0, \tau)) + 2\gamma (s G(s, \tau) - g(0, \tau)) \\ + \omega^2 G(s, \tau) = e^{-s\tau} \end{aligned}$$

Take $g(t, \tau)$ to obey same initial conditions as $x(t)$.

$$\text{So } g(0, \tau) = \dot{g}(0, \tau) = 0,$$

$$(s^2 + 2\gamma s + \omega^2) G(s, \tau) = e^{-s\tau}$$

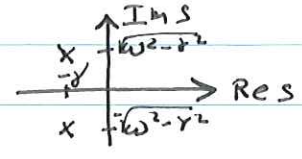
$$G(s, \tau) = \frac{e^{-s\tau}}{s^2 + 2\gamma s + \omega^2} = \frac{e^{-s\tau}}{(s - \lambda_+)(s - \lambda_-)}$$

$$\text{where } \lambda_{\pm} = \frac{1}{2}(-2\gamma \pm \sqrt{4\gamma^2 - 4\omega^2}) = -\gamma \pm \sqrt{\gamma^2 - \omega^2}$$

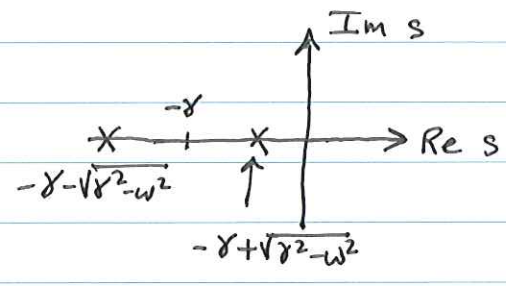
$G(s, \tau)$ has ^{simple} poles at $s = \lambda_+$ and $s = \lambda_-$.

(i) case $\omega > \gamma \rightarrow \pm \sqrt{\gamma^2 - \omega^2} = \pm i \sqrt{\omega^2 - \gamma^2}$ is imaginary.

\rightarrow poles at $s = -\gamma \pm i \sqrt{\omega^2 - \gamma^2}$



(ii) case $\omega < \gamma \rightarrow \pm \sqrt{\gamma^2 - \omega^2}$ is real



poles are always in the $\text{Re}(s) < 0$ plane. Can take Bromwich contour $\int_{-i\infty}^{i\infty} ds$.

$$g(t, \tau) = \frac{1}{2\pi i} \int_{-i\infty}^{i\infty} ds e^{st} G(s, \tau)$$

$$= \frac{1}{2\pi i} \int_{-i\infty}^{i\infty} ds \frac{e^{(t-\tau)s}}{(s-\lambda_+)(s-\lambda_-)}$$

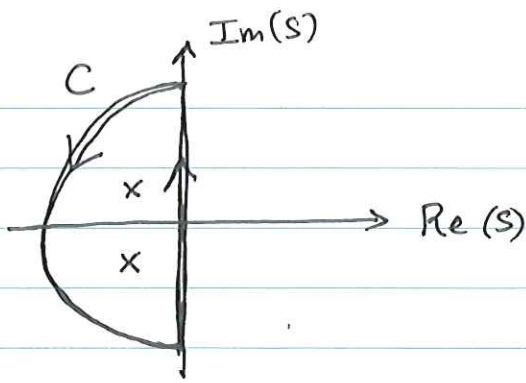
How do we want to close the contour? Want the integrand to go to zero when we close the contour.

case 1: $t > \tau$.

$e^{(t-\tau)s} \rightarrow 0$ for $\text{Re}(s) \rightarrow -\infty$

$e^{(t-\tau)s} \rightarrow \infty$ for $\text{Re}(s) \rightarrow +\infty$

close the contour in the $\text{Re}(s) < 0$ plane.



$$\begin{aligned}
 g(t, \tau) &= \frac{1}{2\pi i} 2\pi i \left\{ \text{Res}(\lambda_+) + \text{Res}(\lambda_-) \right\} \\
 &= \lim_{s \rightarrow \lambda_+} (s - \lambda_+) \frac{e^{(t-\tau)s}}{(s - \lambda_+)(s - \lambda_-)} + \text{same} (\lambda_+ \leftrightarrow \lambda_-) \\
 &= \frac{e^{(t-\tau)\lambda_+}}{\lambda_+ - \lambda_-} + \frac{e^{(t-\tau)\lambda_-}}{\lambda_- - \lambda_+} \\
 &= \frac{1}{2\sqrt{\gamma^2 - \omega^2}} \left[e^{(t-\tau)\lambda_+} - e^{(t-\tau)\lambda_-} \right]
 \end{aligned}$$

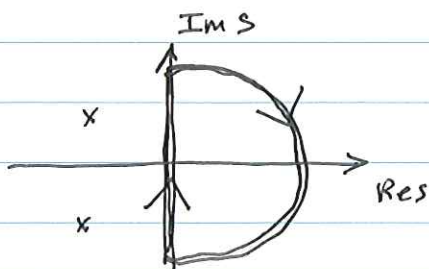
Let's work in the limit $\gamma \ll \omega$ for simplicity:

$$\begin{aligned}
 g(t, \tau) &= \frac{1}{2i\omega} \left[e^{-\gamma(t-\tau) + i\omega(t-\tau)} - e^{-\gamma(t-\tau) - i\omega(t-\tau)} \right] \\
 &= \frac{1}{\omega} e^{-\gamma(t-\tau)} \sin(\omega(t-\tau))
 \end{aligned}$$

Case 2: $t < \tau$.

$$\begin{aligned}
 e^{(t-\tau)s} &\rightarrow \infty \text{ for } \text{Re}(s) \rightarrow -\infty \\
 e^{(t-\tau)s} &\rightarrow 0 \text{ for } \text{Re}(s) \rightarrow +\infty
 \end{aligned}$$

close contour in the $\text{Re}(s) \gg 0$ plane



~~Res~~ No enclosed poles.
So $g(t, \tau) = 0$

So we have:
$$g(t, \tau) = \begin{cases} 0 & t < \tau \\ \frac{1}{\omega} e^{-\gamma(t-\tau)} \sin(\omega(t-\tau)) & t > \tau \end{cases}$$

$$= \theta(t-\tau) \frac{1}{\omega} e^{-\gamma(t-\tau)} \sin(\omega(t-\tau))$$

Note: 1. response only depends on $t-\tau$ (time ~~at~~ elapsed after the impulse), not the absolute time t .

2. $g(t, \tau) = 0$ for $t < \tau \rightarrow$ causality. There is no response before the impulse occurs.

For any forcing function $F(t)$, the solution for $x(t)$ is:

$$x(t) = \int_0^{\infty} d\tau \frac{1}{\omega} \theta(t-\tau) e^{-\gamma(t-\tau)} \sin(\omega(t-\tau)) F(\tau)$$

$$= \int_0^t d\tau \frac{1}{\omega} e^{-\gamma(t-\tau)} \sin(\omega(t-\tau)) F(\tau)$$

Solution ~~is~~ $x(t)$ only depends on times $\tau < t$, not on times $\tau > t$ (causality)