



Higgs

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1 The Georgi-Glashow SU(2) Model

The Georgi-Glashow model was proposed as a theory of the weak interactions in 1972. It doesn't quite work, as you will see, but many key features of the electroweak theory are present. The Lagrangian for the Georgi-Glashow model is:

$$\mathcal{L}_{\text{GG}} = -\frac{1}{2}\text{Tr}(F_{\mu\nu}F^{\mu\nu}) + \text{Tr}[(D_\mu\phi)(D^\mu\phi)] + \mu^2\text{Tr}\phi^2 - \lambda(\text{Tr}\phi^2)^2 \quad (1)$$

The gauge group is $SU(2)$, and the scalar field $\phi = \phi^i T^i$ is a triplet (i.e. adjoint) whose components ϕ^i are real and

$$T^a = \frac{1}{2}\sigma^a \quad (2)$$

- a) What values of ϕ minimize the potential?
- b) By $SU(2)$ symmetry you can choose the vacuum to be $\langle\phi^1\rangle = \langle\phi^2\rangle = 0$, $\langle\phi^3\rangle = v$.

Show that this choice of vacuum “breaks” two of the generators of $SU(2)$

$$e^{i\alpha_{1,2}T^{1,2}}\langle\phi\rangle e^{-i\alpha_{1,2}T^{1,2}} \neq \langle\phi\rangle \quad (3)$$

and leaves one generator unbroken

$$e^{i\alpha_3 T^3}\langle\phi\rangle e^{-i\alpha_3 T^3} = \langle\phi\rangle \quad (4)$$

by using the infinitesimal versions of equations (3) and (4).

- c) Show that you can make a gauge choice (known as unitary gauge) so that $\phi^1 = \phi^2 = 0$, $\phi^3 = v + h$.
- d) In unitary gauge, show that two of the three gauge bosons get a mass $M_A^2 = g^2 v^2$, and that one of the gauge bosons remains massless.

The Georgi-Glashow model can accommodate the massive W^\pm bosons as well as the photon, but it has no Z boson. The Z boson is necessary to explain weak neutral current interactions, which were first observed in elastic electron-neutrino scattering.

2 *The Higgs Sector*

This problem is essential for this course. The Lagrangian for the Higgs sector in the Standard Model is given by

$$\mathcal{L}_{\text{Higgs}} = (D_\mu H)^\dagger (D^\mu H) - V(H)$$

where H is an $SU(2)_L$ doublet of scalar fields and

$$V(H) = -\mu^2 H^\dagger H + \lambda (H^\dagger H)^2$$

is the $SU(2)_L \times U(1)_Y$ invariant potential.

Consider a constant value of H for which this potential is minimized. By $SU(2)$ symmetry we can always choose this value to be

$$\begin{pmatrix} 0 \\ \frac{v}{\sqrt{2}} \end{pmatrix} \quad (5)$$

where v is a constant. Recall that the $U(1)_Y$ charge of H is $\frac{1}{2}$.

a) Show that the generator of $SU(2)_L \times U(1)_Y$ that leaves the vacuum invariant is $T_3 + Y$.

Next consider small fluctuations of the Higgs field around this vacuum in the unitary gauge

$$H(x) = \begin{pmatrix} 0 \\ \frac{v+h(x)}{\sqrt{2}} \end{pmatrix}.$$

b) Show that the potential then takes on the form

$$V(h) = (\lambda v^2)h^2 + \lambda v h^3 + \frac{\lambda}{4}h^4.$$

c) Also show that the kinetic term for the Higgs boson in the Lagrangian becomes

$$\begin{aligned} \mathcal{L}_{\text{Kinetic}} = & \frac{1}{2}(\partial_\mu h)(\partial^\mu h) + \frac{1}{8}(g'B_\mu - gX_\mu^3)(g'B^\mu - gX^{3\mu})(v+h)^2 \\ & + \frac{1}{8}(gX_\mu^1 - igX_\mu^2)(gX^{1\mu} + igX^{2\mu})(v+h)^2 \end{aligned}$$

which shows that some of the vector bosons gain mass terms.

d) If we perform a field redefinition, we want the kinetic terms for the vector bosons to remain canonical normalized. Show that correct mass-eigenstate for the neutral particle above is

$$Z_\mu = \cos \theta_W X_\mu^3 - \sin \theta_W B_\mu$$

where

$$\cos \theta_W = \frac{g}{\sqrt{(g')^2 + g^2}}$$

is the Weinberg angle.

- e) Verify that under a gauge transformation involving $U(1)_{EM}$ ($U(1)_{EM}$ is the unbroken subgroup of $SU(2)_L \times U(1)_Y$), X_μ^1 and X_μ^2 transform into each other, and thus

$$W^{\pm\mu} = \frac{1}{\sqrt{2}} (X_\mu^1 \mp iX_\mu^2)$$

represent charged particles under $U(1)_{EM}$.

- f) Verify that the masses of the three massive bosons are given by

$$m_W^2 = \frac{1}{4}g^2v^2 \quad \text{and} \quad m_Z^2 = \frac{1}{4}v^2((g')^2 + g^2)$$

- g) What is the expression for the vector field A_μ (given in terms of g, g', X_μ^a and B_μ) that remains massless? What is it called?
- h) In general, the covariant derivative of field with $SU(2)_L \times U(1)_Y$ charge that is an $SU(2)$ doublet is given by

$$D_\mu = \partial_\mu - igX_\mu^a T^a - ig'Y B_\mu \quad (6)$$

where

$$T^a = \frac{1}{2}\sigma^a \quad (7)$$

Please rewrite this covariant derivative in terms of W_μ^\pm, Z_μ, A_μ and obtain the following form:

$$D_\mu = \partial_\mu - i\frac{g}{\sqrt{2}}(W_\mu^+ T^+ + W_\mu^- T^-) - i\frac{g}{\cos\theta_W}Z_\mu(T^3 - Q \sin^2\theta_W) - ieA_\mu Q \quad (8)$$

where we also define

$$T^\pm \equiv T^1 \pm iT^2 \quad (9)$$

and

$$Q \equiv T_3 + Y, \quad e \equiv \frac{gg'}{\sqrt{g^2 + g'^2}}$$