

Rotational dynamics

①

fundamental equations to describe
dynamics

translations

$$m\vec{a} = \sum \vec{F}$$

equilibrium

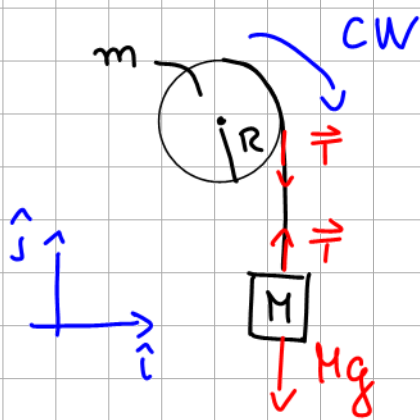
$$\sum \vec{F} = 0$$

rotations
(about fixed axis)

$$I\vec{\alpha} = \sum \vec{\tau}^*)$$

$$\sum \vec{\tau} = 0$$

example of rotational dynamics: pulling a pulley



$$\vec{\tau} = \vec{r} \times \vec{F} = \vec{r} \times \vec{T} = -RT\hat{k} = I\vec{\alpha}$$

$$\text{CO } \alpha_z = -\frac{RT}{I} = -\frac{2T}{mR}$$

$I = \frac{m}{2}R^2$ for disk (Table 8.2)

Newton-2 for crate: $F_y = -Mg + T = Ma_y$
 $\Leftrightarrow T = M(a_y + g)$

insert in α_z equation (to eliminate tension):

$$\text{CO } \alpha_z = -\frac{2M}{mR} (a_y + g)$$

This is a relation between α_z and a_y . Another relation between these two (yet unknown) quantities is

$$a_y = a_t = \alpha_z R$$
$$\Leftrightarrow \alpha_z = \frac{a_y}{R}$$
$$= -\frac{2M}{mR} (a_y + g)$$

(a_y = tangential acceleration of pulley if rope doesn't become slack)

* One can derive this equation from Newton-2 for any object with moment of inertia I .

(2)

$$\Rightarrow a_y + \frac{2M}{m} a_y = -\frac{2M}{m} g$$

$$\Rightarrow a_y \left(\frac{m+2M}{m} \right) = -\frac{2M}{m} g$$

$$\Rightarrow a_y = -\frac{2Mg}{m+2M} = -\frac{Mg}{M + \frac{m}{2}} \underset{m \ll M}{\approx} -g$$

$$\text{CO } \alpha_z = \frac{a_y}{R} = -\frac{Mg}{R(M + \frac{m}{2})} \underset{m \ll M}{\approx} -\frac{g}{R}$$

(simple free fall if $m \ll M$)

$$T = M(a_y + g) = Mg \left(1 - \frac{2M}{m+2M} \right) = \frac{Mg}{\frac{2M}{m} + 1} \underset{m \ll M}{\approx} \frac{Mg}{\frac{2M}{m}} = \frac{mg}{2}$$

Both a_y and α_z are constant. Hence, we can infer

(angular) velocity and (angular) displacement using

$$\omega_z(t) = \omega_0 + \alpha_z t, \quad \theta(t) = \theta_0 + \frac{\alpha_z}{2} t^2 \underset{m \ll M}{\approx} \theta_0 - \frac{gt^2}{2R}$$

$$v_y(t) = v_0 + a_y t, \quad y(t) = y_0 + \frac{a_y}{2} t^2 \underset{m \ll M}{\approx} y_0 - \frac{gt^2}{2}$$