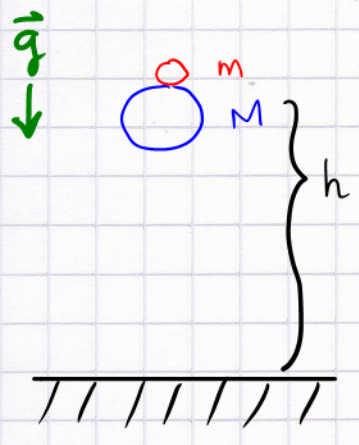


"Astrobounce"

A pair of balls, m sits on top of M , is dropped from height h . What happens?



A sequence of "events"

- 1) M moves from h to ground (= Bottom)
- m " " " " almost ground (= Bottom)

$$v_B = V_B = \sqrt{2gh}$$

speeds (> 0)

convert

$$Mgh \rightarrow \frac{1}{2} M V_B^2$$

$$mgh \rightarrow \frac{1}{2} m v_B^2$$

- 2) M (at the bottom) compresses, converts $KE = \frac{1}{2} M V_f^2 \rightarrow$ spring-like energy \rightarrow reverses direction while m is still moving down \rightarrow collision

Detailed analysis: make a \hat{j} choice, e.g., \uparrow , then:

$$M V_B - m v_B = M V_f + m v_f$$

\uparrow reversed, going up

\nwarrow coming down!

speeds at bottom \rightarrow

$$V_B = v_B = \sqrt{2gh} \quad (> 0)$$

V_f, v_f are velocities (with sign) after collision

$$(M-m)\sqrt{2gh} = M V_f + m v_f \rightarrow V_f = (M-m)\sqrt{2gh} - m v_f$$

- A) One equation, two unknowns (V_f, v_f) cannot solve
- B) Why was momentum conservation valid? The system is NOT force free! Gravity + normal force are external forces acting on M , gravity acts on m .

Answer: collision takes place over short times ($\Delta t \sim \text{msec}$) External forces cannot change P^{tot} by much over Δt

- C) Energy conservation during (perfectly) elastic collision:

$$\frac{1}{2} M (2gh) + \frac{1}{2} m (2gh) = \frac{1}{2} M V_f^2 + \frac{1}{2} m v_f^2$$

eliminate V_f
+ solve for v_f
TEDIOUS! doable!

Simplified analysis for the M-ball + m-ball collision: ②

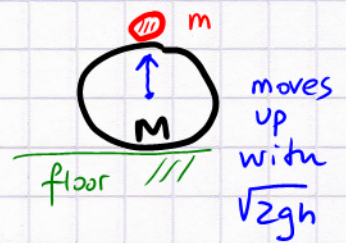
go to the reference frame moving with M:

"back of envelope"

in this frame: m comes down with $2\sqrt{2gh}$:

$$M \cdot 0 - m(2\sqrt{2gh}) = M \tilde{V}_f + m \tilde{v}_f$$

$$\therefore \tilde{V}_f = -\frac{m}{M}(2\sqrt{2gh}) - \frac{m}{M}\tilde{v}_f$$



In this frame M will see a (small) recoil (downward velocity, $\frac{m}{M}$ factor \rightarrow small)

$M \gg m$ analysis: in this frame, m "bounces" from a wall \rightarrow reverses its velocity $\rightarrow \tilde{v}_f \approx 2\sqrt{2gh}$ up

With respect to the floor: $v_f = \tilde{v}_f + v_{\text{frame}} = 3\sqrt{2gh}$
(absolute space) \uparrow frame \uparrow M moves up with $\sqrt{2gh}$!

\therefore In $M \gg m$ limit: top ball goes from $\sqrt{2gh}$ down $\rightarrow 3\sqrt{2gh}$ up

Now convert KE \rightarrow PE for top ball:

$$mgh = \frac{1}{2} m v_f^2 = \frac{1}{2} m \cdot 9 \cdot 2gh$$

$$H = 9h \quad \text{nine-fold original height!}$$

Physics message: • collisions allow transfer of momentum

• this also implies transfer of KE

• big ball has a lot more energy: Mgh vs mgh ($M \gg m$)

• transfers a big part of Mgh (almost comes to rest \rightarrow BOTTOM)

• top ball gains this energy; theoretically could reach up to $H = 9h$.

Appendix

Do the full calculation in the "normal" frame and then look at $M \gg m$.

A-2

$$(M-m)\sqrt{2gh} = MV_f + mv_f \quad \textcircled{1} \quad \text{total momentum conservation}$$

$$\frac{1}{2}M(2gh) + \frac{1}{2}m(2gh) = \frac{1}{2}MV_f^2 + \frac{1}{2}mv_f^2 \quad \textcircled{2} \quad \text{energy conservation during elastic collision}$$

How do we solve for v_f and V_f ?

Use ① to eliminate V_f in ②:

$$V_f = \left(1 - \frac{m}{M}\right)\sqrt{2gh} - \frac{m}{M}v_f$$

$$\textcircled{2}: 2gh(M+m) = MV_f^2 + mv_f^2$$

$$\therefore 2gh(M+m) = M \left(\left(1 - \frac{m}{M}\right)\sqrt{2gh} - \frac{m}{M}v_f \right)^2 + mv_f^2$$

Sort out the mess \rightarrow quadratic eqn in v_f ...

$$2gh \frac{m+3M}{m+M} = v_f^2 - 2\sqrt{2gh} \frac{M-m}{M+m} v_f$$

$$v_f = \sqrt{2gh} \frac{M-m}{M+m} \pm \sqrt{2gh \left(\frac{M-m}{M+m}\right)^2 + 2gh \frac{3M+m}{M+m}}$$

$$= \sqrt{2gh} \left[\underbrace{\frac{M-m}{M+m}}_{\approx 1 \text{ for } M \gg m} + \sqrt{\frac{(M-m)^2 + (3M+m)(M+m)}{(M+m)^2}} \right]$$

↑ $v_B!$ ↑ increase!

$$= \frac{v_B}{M+m} \left[M-m + \sqrt{\frac{M^2 - 2mM + m^2 + 3M^2 + 4mM + m^2}{4M^2 + 2mM + 2m^2}} \right]$$

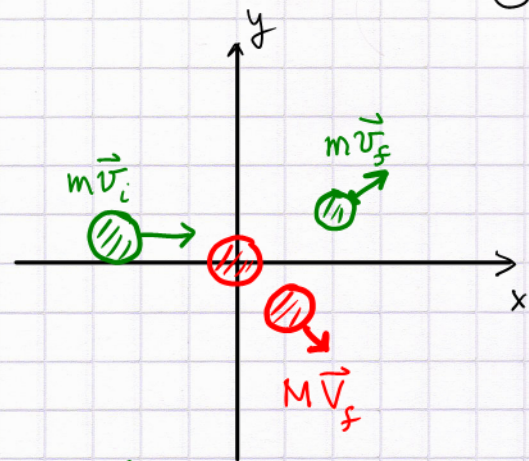
$$\begin{aligned} & \frac{M+2M}{M} v_B \\ & \rightarrow 3v_B \\ & M \gg m \end{aligned}$$

Elastic Collisions in 2D

hockey pucks colliding off-center

initially: $M\vec{V}_i = \vec{0}$

$$m\vec{v}_i = m v \hat{i}$$



finally: $M\vec{V}_f$, $m\vec{v}_f \rightarrow m v_{fx} \hat{i} + m v_{fy} \hat{j}$

$$\rightarrow M V_{fx} \hat{i} + M V_{fy} \hat{j}$$

Momentum conservation: vector eqn \rightarrow 2 eqs

$$m v_{ix} + 0 = m v_{fx} + M V_{fx} ; 0 = m v_{fy} + M V_{fy}$$

Energy conservation: one scalar eqn

$$\frac{1}{2} m v_{ix}^2 = \frac{1}{2} m (v_{fx}^2 + v_{fy}^2) + \frac{1}{2} M (V_{fx}^2 + V_{fy}^2)$$

Directional info: $v_{fx} = v_f \cos \theta$, $v_{fy} = v_f \sin \theta$

$$V_{fx} = V_f \cos \varphi, V_{fy} = -V_f \sin \varphi$$

Q: how would we show that $(\theta + \varphi) \approx \frac{\pi}{2}$ as demonstrated in the video?
for peripheral collisions \leftarrow

$m=M$ case $\rightarrow v \sin \theta = -V \sin \varphi$, $v_0 = v \cos \theta + V \cos \varphi$,
 $v_0^2 = v^2 + V^2$;

Note: we have 3 equations. For given input momenta there are 4 unknowns: (v, V, θ, φ) or (p_x, p_y, P_x, P_y)

\therefore we cannot determine the outcome uniquely!

We can discuss special cases.

1) Equal energy sharing $v = V$; since $m = M$ by assumption

$$\frac{1}{2} m v^2 = \frac{1}{2} M V^2$$

What follows?

$$v_0^2 = v^2 + V^2 = 2v^2 \quad \therefore v = \frac{v_0}{\sqrt{2}}$$

$$v \sin \theta = -V \sin \varphi \quad \therefore v \sin \theta = -v \sin \varphi \quad \left| \begin{array}{l} \text{but} \\ v \neq 0 \end{array} \right.$$

$$\therefore \theta = -\varphi$$

$$\begin{aligned} \text{Now use: } v_0 &= v \cos \theta + v \cos \varphi = v (\cos \theta + \cos(-\theta)) \\ &= \sqrt{2} v = 2v \cos \theta \end{aligned}$$

$$\therefore \cos \theta = \frac{1}{2} \sqrt{2} \quad \therefore \theta = \frac{\pi}{4}, \quad \varphi = -\frac{\pi}{4}$$

$$\text{Note: } \theta + |\varphi| = \frac{\pi}{2}$$

Equal-mass pucks collide :
 m hits stationary $M \rightarrow$
 for both to come out with equal
 energy they emerge as

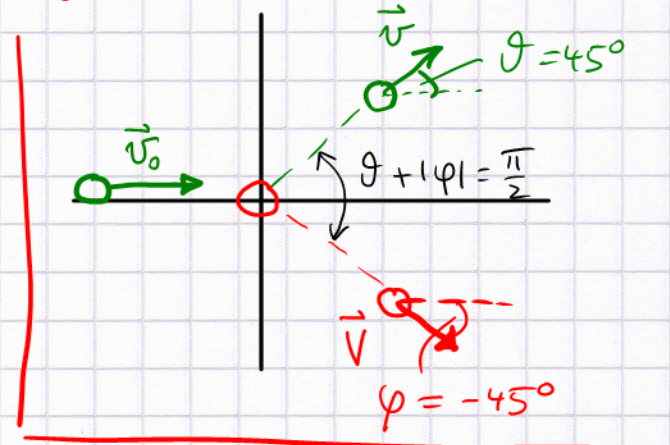
2) A grazing collision:

small θ (little deflection)

expect $v \approx v_0, V \ll v_0$.

(from $v_0^2 = v^2 + V^2$)

$$p_y = m v \sin \theta = -p_y = -M V \sin \varphi$$



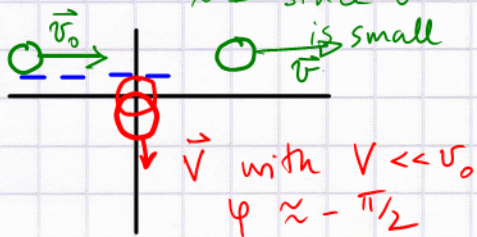
if $\sin \theta$ is small & v is large \therefore if V is small, $\sin \varphi$ is large

$$v_0 = v \underbrace{\cos \theta}_{\approx 1 \text{ since } \theta \text{ is small}} + V \cos \varphi \quad \therefore \cos \varphi = \frac{v_0 - v}{V} \quad \frac{\text{small}}{\text{small}} ?$$

How can we make the argument tight?

$$v_0^2 - v^2 = (v_0 - v)(v_0 + v) = V^2$$

$$\therefore v_0 - v \approx \frac{V^2}{2v_0} \quad \therefore \cos \varphi \approx \frac{V}{2v_0} \quad \begin{array}{l} \text{really} \\ \text{small} \end{array} \rightarrow \varphi = -\frac{\pi}{2}$$



\vec{V} with $V \ll v_0$
 $\varphi \approx -\frac{\pi}{2}$ again: $\theta + |\varphi| \approx \frac{\pi}{2}$.