

EM Waves

C24W10

G23.1-2

From Faraday's law (Lenz) we saw that an EMF is generated in a conducting loop of fixed size and orientation, if the strength of the magnetic field varies with time

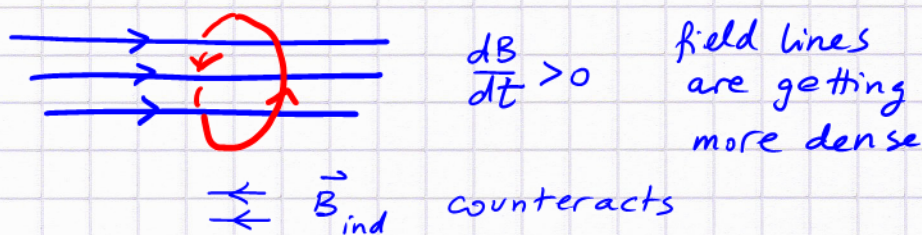
$$\Phi_M = B A \cos \theta : \quad \frac{dB}{dt} \rightarrow \frac{d\Phi_M}{dt} \rightarrow \mathcal{E}$$

Lenz: There is an induced current in the loop to generate an induced \vec{B} field to counteract $\frac{d\Phi_M}{dt}$.

The EMF \mathcal{E} responsible for the induced current is associated with an \vec{E} field (otherwise no induced current) which is very different from previously discussed \vec{E} fields:

Previously: \vec{E} starts at $q > 0$ charge and ends at $q < 0$ charge

Here: \vec{E} goes around in a loop without beginning and end.



When the time variation is very rapid (high frequency begins at $\sim 100 \text{ kHz} = 10^5 \text{ Hz}$) it was found ~ 1900 by Hertz that this form of induction can happen over distances.

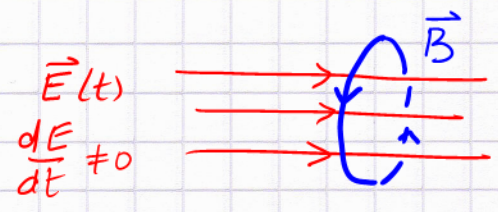
LC oscillator/receiver circuits can communicate at such and higher frequencies \rightarrow the induced \vec{E} field can exist in air (vacuum) without a conducting metal loop.

To fully understand the propagation of electric/magnetic fields one also had to figure out the following symmetry:

A time-varying \vec{E} field surrounds itself with a circular \vec{B} field.

Note: \vec{E} generated from Φ_M :
 $\vec{E} \perp \vec{B}$

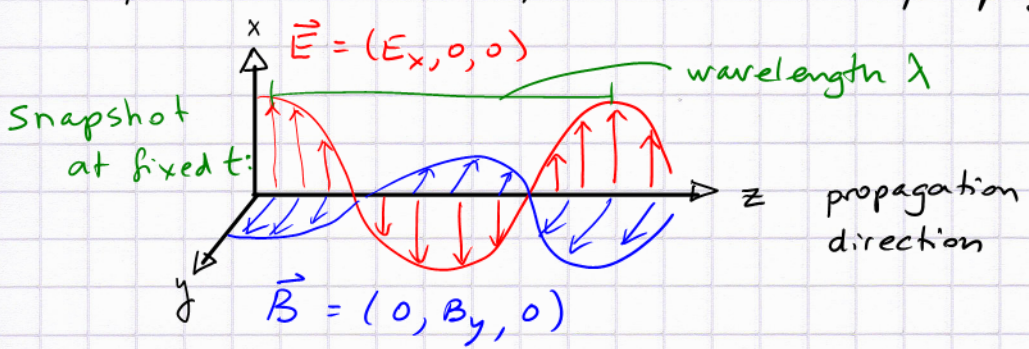
and \vec{B} generated from Φ_E :
 $\vec{B} \perp \vec{E}$



time varying Φ_E through virtual loop $\rightarrow \vec{B}$ is generated

The symmetry: $\frac{d}{dt} \Phi_M \rightarrow \vec{E}$ Faraday law
 $\frac{d}{dt} \Phi_E \rightarrow \vec{B}$ (generalized Ampère law)

forms the basis for EM wave propagation through free space see Fig. 23.3



\vec{E} and \vec{B} are in phase

$\vec{E} \times \vec{B} \sim \hat{k}$
propagation direction

when such a wave encounters charges (e.g., bound electrons in glass)

$\rightarrow \vec{E}$ sets them in oscillatory motion

$\vec{E} \perp \hat{k}$ and $\vec{B} \perp \hat{k}$
 \Rightarrow transverse wave

When a charge oscillates rapidly

\rightarrow it gives off EM radiation with the oscillation frequency

antenna \rightarrow radiowaves, microwaves

Characterization of an EM wave

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1) shown above is a special form:

\vec{E}, \vec{B} oscillate in fixed, perpendicular directions

→ linearly polarized EM wave

(lasers can produce such waves)

2) the wave is a very long train of regular oscillations

practical durations:

nsec - laser

psec - laser

fsec → "short" pulse

atto second → ultra-short pulse

(Paul Corkum in Ottawa: world leader!)

3) many conventional light sources

- discharge lamp* (atomic transitions from excited to lower levels)

↳ nsec bursts of single frequency/wavelength light

* contains a gas, e.g., H_2

apply a high voltage (dangerous → LAB# ? !)

electron breakdown current → dissociates $H_2 \rightarrow H+H$
+ excites H atoms

$H(2p \rightarrow 1s)$ is strong ultraviolet → filtered by glass

$H(3d \rightarrow 2p)$

$H(4f \rightarrow 3d, \text{ or } 4p \rightarrow 3s)$

...

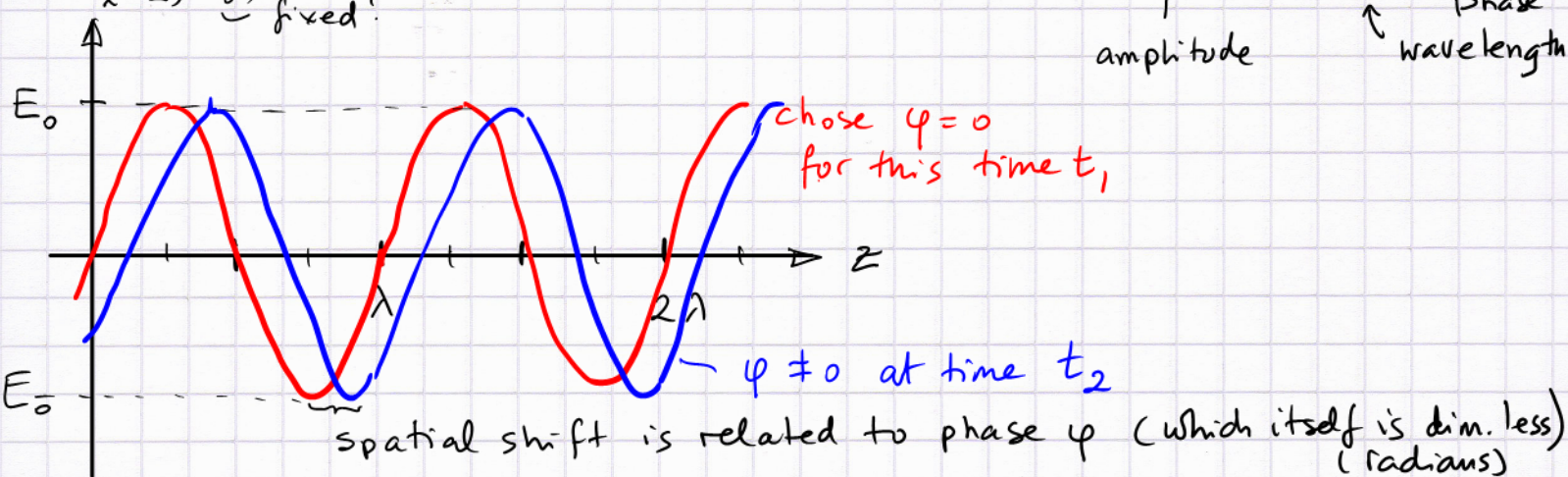
emission lines!

→ see Fig 29.7 p. 992!

Light from these sources → interference → wave model is needed!
(Michelson interferometer)

Two ways to look at the \vec{E}_x field:

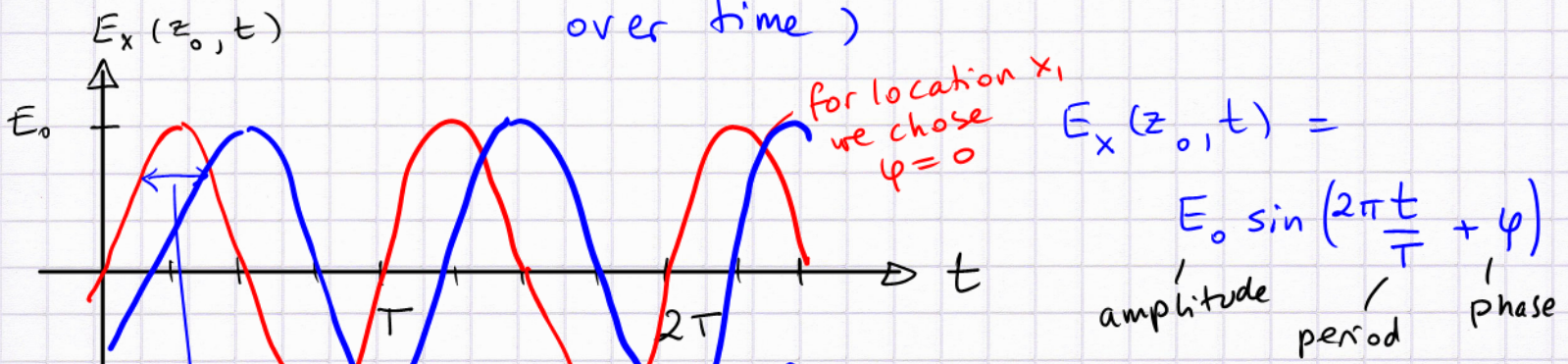
(1) spatial (fixed-time snapshot) $E_x(t_0) = E_0 \sin(\frac{2\pi z}{\lambda} + \varphi)$
 $E_x(z, t_0)$ fixed!
 $t = t_0$
amplitude
wavelength



For visible-light EM waves: $400 \text{ nm} < \lambda < 700 \text{ nm}$

Short WL \rightarrow UV
Long WL \rightarrow infrared (heat)

(2) temporal (fix location x_0 , look how $E_x(z_0, t)$ changes over time)



$$E_x(z_0, t) = E_0 \sin\left(\frac{2\pi t}{T} + \varphi\right)$$

amplitude period phase

Frequencies for visible light: hundreds of THz
Mega, Giga, Tera = 10^{12}

On dimensional grounds: the spatial (λ) and temporal (T , or $f = \frac{1}{T}$) properties are related

by the propagation speed $c = 3.0 \times 10^8 \text{ m/s}$

$$\lambda \cdot f = c$$