

Waves

C25W10
G 12.1-2

In addition to : (point) particles, extended objects (described by centre-of-mass motion + rotation), fields (\vec{E}, \vec{B}) we were confronted with electromagnetic waves, i.e., (\vec{E}, \vec{B}) fields that propagate in space-time.

Waves arise in simpler contexts : water waves, acoustic (sound) waves, waves on a string (standing vs travelling).

What exactly is a wave ?

→ vibrating string (guitar, violin, piano)
some elastic mass is moving up/down

→ water wave : water droplets are sloshing up/down

→ acoustic wave : air is compressed/decompressed along the direction of travel

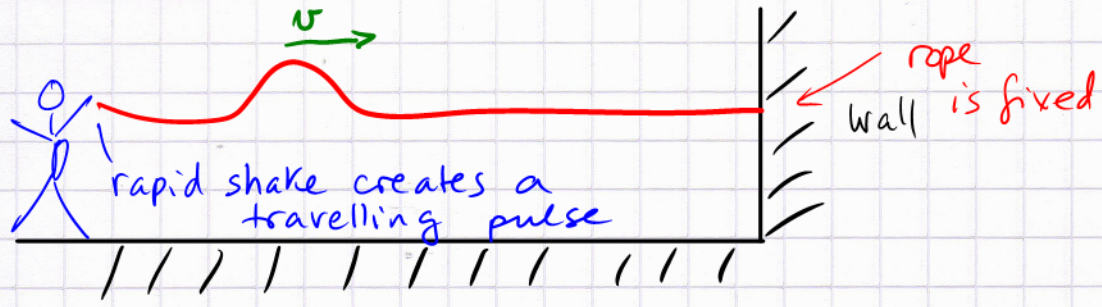
→ waves can travel, but don't have to (→ string)
↓
fixed ends

→ Some medium oscillates, oscillation can be transverse or longitudinal

no medium → no sound propagation (demo!)

Electromagnetic waves are special : they propagate in vacuum!
(induction makes this possible)

Start with travelling transverse wave pulse on a rope



- Questions:
- what controls the propagation speed? (1)
 - what happens to the pulse at the fixed end? (2)
 - does the shape of the pulse change as it travels down the rope and back? (3)
 - how do the particles move that make up the rope? (4)
 - wave pulse vs sinusoid ? (5)

(1) → what can we change about the rope or string?

(A) thick vs thin string (piano strings? / or class. guitar)
 $\mu = M/L = \frac{\text{mass}}{\text{length}}$

(B) tensioning of string → tuning
 tension force F_T is adjusted for fixed length string

Dimensional analysis: $\frac{F_T}{\mu} = \text{accel.} \times \text{length} \rightarrow (\text{velocity})^2$

wave propagation speed $v = \sqrt{\frac{F_T}{\mu}}$ Eq.(12.9) transverse pulses more fast or slow according to this equation

Another example: shallow water (beach) → $\text{depth} \ll \text{wavelength}$

What can the propagation speed depend on?

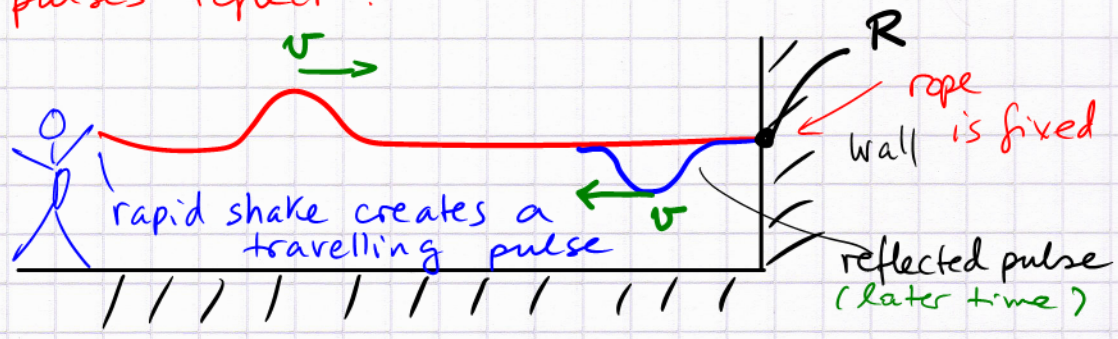
- (A) depth of water, d
- (B) gravitational acceleration g

$v = \sqrt{gd}$ is the simplest formula on dimensional grounds

Q: does it work?

Q2: How do pulses reflect?

see Fig. 12.17 in G. compare with 12.20



At point R the rope is fixed. Pulses reverse at R by flipping their shape vertically (Observation)

Q3 Apart from flipping the shape remains approximately the same.

However: water surface waves tend to dissipate (luckily for canoeists, when one stays away from a passing motor boat ...)

Yet, there exist water waves that propagate over huge distances (tsunami; soliton waves)

→ bottom line: it's complicated, in general

Q4: Who moves, and in which direction?

(1) small-amplitude waves on a string:

- mass segments Δm move up and down
- yet the pulse shape travels left \leftrightarrow right

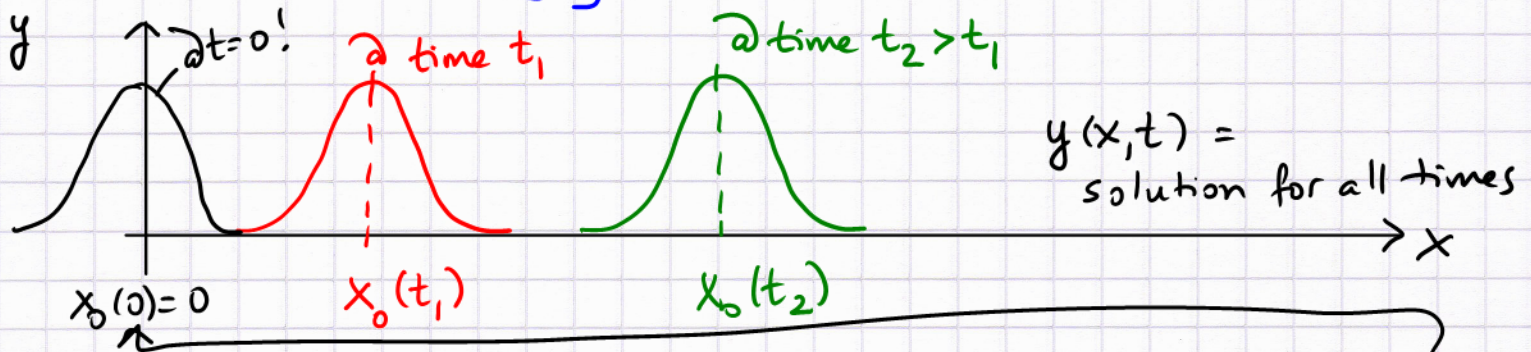
(2) water surface waves: water droplets carry out a circular motion \rightarrow up/down and left/right

- they do not propagate over larger distances

Travelling wave pulse

C25W12p.2

Q: how do we describe a "bump" travelling without the bump shape changing over time?



$$x_0(t) = \text{average position of bump} = v_w t \quad (=0 \text{ at } t=0)$$

$$v_w = c = \text{wave propagation speed}$$

At $t=0$ simple: $y(x,0) = f(x)$ ← bump shape

Examples: $f(x) = \exp\left(-\left(\frac{x}{w}\right)^2\right)$ $w \sim$ width of bump

could also define triangle shape:



Question: how do we make $f(x)$ "move" in time?

Answer: replace $f(x) \rightarrow f(x-ct)$

Why? The peak in f occurs for argument = 0,

$$\text{argument} = x - ct = 0 \quad \therefore x = ct$$

yields x_0

Example:

$$f(x-ct) = \exp\left[-\frac{(x-ct)^2}{w_0^2}\right] \quad \text{travels to the right (for } c > 0)$$

Sound waves do not propagate in vacuum

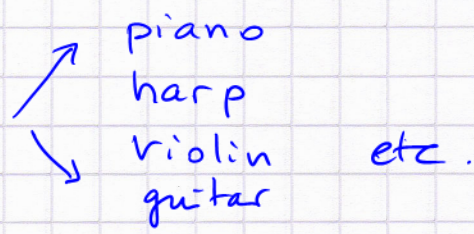
they are longitudinal (along travel direction)

compressions + decompressions (rarefactions)

caused by the movement of

1) speaker membrane

2) instrument 'sound board'



the strings are attached to a large-area board + make it vibrate

3) sound chords co-operate with oral cavity

4) wind instruments (pipe, flute, horn, etc.)

lips/mouthpiece create vibrations in air column of certain length

Longitudinal waves: particles oscillate in the direction of the wave propagation, but do not travel far

Common to all of this:

propagation of wave (energy transport)

has little to do with the motion of the medium!

⇒ wave = new concept

it is more than the motion of the medium

Wave pulse vs wave train (Q5)

(1) Travelling wave pulse is characterized by some shape function D ; at fixed time $t_0 = 0$:

$$D(x) = \text{graph of pulse centered at } x_0 \text{ with width } w, \text{ e.g., } A e^{-\frac{(x-x_0)^2}{w^2}}$$

Gaussian: $x \mapsto \exp(-x^2/w^2)$ $w \rightarrow$ width parameter

pulse is centered on location x_0

$A =$ amplitude ($\exp(-0^2/w^2) = 1$)

For any time t : $A \exp\left(-\left[\frac{x - (x_0 + vt)}{w}\right]^2\right)$

a) $t = 0$ recover the above shape

a) $t > 0$; centre location travels to $x_0 + vt$

The width parameter w is chosen as a constant
 \rightarrow pulse preserves its shape

(2) An infinitely long wave train (Eq. 12.1 in G 12.2)

$$D(x, t) = A \sin\left(2\pi f t - \frac{2\pi x}{\lambda}\right)$$

is characterized by the same periodic (sine) function

in space (fix t , and observe) \rightarrow period = wavelength λ

and in time (fix x , observe) \rightarrow period = T , frequency $f = \frac{1}{T}$

We can simplify to $A \sin(\omega t - kx)$ $\omega = 2\pi f$ circular frequency

Note: $\lambda \cdot f = c \leftarrow$ propagation speed ($=v$)
(Eq. 12.2)

$k = \frac{2\pi}{\lambda} =$ wavenumber

Later in math/phys we learn: a finite wave pulse can be decomposed into a mix (superposition) of sinusoidal wave trains (Fourier analysis). Arbitrary pulse = mix of many fixed frequency/wavelength periodic wave trains.